Orientation and effects of channel H₂O and CO₂ in cordierite

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Abstract
Channel CO₂ and H₂O has been re-introduced into a channel evacuated Mg-cordierite

K₀.₀₁Na₀.₀₃(Mg₁.₉₁Fe₀.₀₉Mn₀.₀₁)Al₃.₉₈Si₅.₀₁O₁₈

from White Well, Australia. At 600°C and pressures up to 6 kbar, a maximum of 3.13 wt.% CO₂ and 2.6 wt.% H₂O re-enter cordierite's channels. The H₂O molecule orients preferentially with its H-H vector parallel to c (optic orientation: c = X, b = Y, a = Z). With increased H₂O content, the refractive indices γ and β increase more strongly than α whereas the a₀ cell edge decreases, b₀ perhaps increases slightly, and c₀ first increases and then levels off. The linear CO₂ molecule orients chiefly parallel to cordierite's a-axis so that, as CO₂ content increased, refractive index T increased more strongly than did α and β. As a result, 2V^ increased with CO₂ content so as to exceed 90°. Simultaneously, c₀ increased, a₀ decreased, and b₀ remained constant. The distortion index Δ for these crystals changed with H₂O and CO₂ content.

Introduction
Cordierites, Naₓ+₁/₂(MgₓFeₓ,Mnₓ)²⁺(Al₄₋ₓBeₓ Si₅)⁴⁺O₁₈ ⋅ n[H₂O,CO₂], consist structurally of six-membered rings of corner-sharing tetrahedra (T₂ in Fig. 1) stacked along the c-axis. These are linked (laterally and vertically) by corner-sharing with other tetrahedra (T₁) to form a framework structure (Gibbs, 1966). The ring-stacking produces channels along the c-axis that pinch to 'bottlenecks' (~2.5 Å in diameter) or swell to large 'cages' whose maximum dimensions, which occur in the plane parallel to (001), are approximately 5.4 Å along b and 6.0 Å along a. In low cordierite, Al concentrates into the two equivalent T₂6 tetrahedra in the Al₂Si₂O₁₈ ring so that, instead of possessing a 6-fold axis, as for the Si₂O₁₈ ring in beryl, the Al₂Si₂O₁₈ ring possesses a 2-fold axis (like that drawn perpendicular to the lightly stippled plane in Fig. 1). This leads to the orthorhombic symmetry (Ccmm, Z = 4, a > b > c) thus far exhibited by cordierites from metamorphic and igneous rocks (Selkregg and Bloss, 1980; Wallace and Wenk, 1980). High-cordierite or indialite, found where shale and sandstone fused near a burning coal seam (Miyashiro and Iiyama, 1954), appears to be hexagonal, the rapid cooling apparently preventing Al from concentrating in T₂6. Such also seems the case for cordierites from meteorites or those synthesized by recrystallization from a glass (Schreyer and Schairer, 1961), by growth in a flux (Lee and Pentecost, 1976), and by hydrothermal methods (Schreyer and Yoder, 1964). Additional heating of such synthetics leads to Si/Al ordering accompanied by twinning along (110) and (310) (Putnis, 1980; Armbruster and Bloss, 1981).

In natural cordierites the structural channels accommodate numerous species of which CO₂ and H₂O can be recognized by IR spectroscopy, and He and Ar by mass spectrometry (Damon and Kulp, 1958; Beltrame et al., 1976). Hydrocarbons and other gases may also be present in small concentrations not revealed in IR spectra. Na⁺ is an important channel occupant that centers in the bottle-necks (Meagher, 1967) and provides charge balance if Be²⁺ substitutes for Al³⁺ (Povondra and Langer, 1971). Small amounts of Fe in the channels may cause the pleochroism of some cordierites (Goldman et al., 1977), but firm data in support of this are not yet at hand.

The degree of variation of the cordierite lattice from hexagonal symmetry is customarily measured
by the distortion index $\Delta$, which Miyashiro (1957) defined (for CuKα radiation) as

$$\Delta = 2\theta_{131} - (2\theta_{311} + 2\theta_{421})/2.$$  

The assumption was then made that $\Delta$ provided a reliable measure of Al/Si ordering for natural cordierites. However, systematic optical and X-ray studies of natural and heated cordierites (Selkregg and Bloss, 1980) has shown this to be unlikely for natural cordierites, excluding indialites. Instead, the $\Delta$ index depended on Na-content, Fe- and Mn-substitution for Mg, and on water content. Although degrees of Si/Al ordering would without doubt influence $\Delta$, cordierites from metamorphic and igneous rocks appear to have cooled so slowly that, thus far, all crystal structure analyses (Gibbs, 1966; Cohen, et al., 1977; Hochella et al., 1979; Wallace and Wenk, 1980) have disclosed nearly perfect Si/Al ordering.

In studies of natural cordierites, the tendency has been to overlook all channel occupants except H2O. However, CO2 also seems important (Suknev et al., 1971; Schreyer et al., 1979; Hörmann et al., 1980) and leads to specific distortions of the cordierite lattice and modifications of its optical properties (Armbruster and Bloss, 1980). The present paper examines the effects of the orientation of H2O and CO2 on the crystal lattice and optical properties of a cordierite with a well ordered Si/Al distribution.

**Experimental**

A Mg-cordierite from White Well, Australia (Pryce, 1973) with well ordered Si/Al distribution (Cohen et al., 1977; Hochella et al., 1979) was crushed to an average grain size of about 200 μm. Transparent, inclusion-free grains were hand picked under a binocular microscope and tempered for one day between 1200°C and 1300°C in a SiC-furnace in order to expel the volatile channel occupants. Electron microprobe and coulombmetrical analyses for H2O and CO2 established this cordierite’s composition to be

$$K_{0.00}Na_{0.05}(Mg_{1.94}Fe_{0.08}Mn_{0.01})Al_{3.97}Si_{4.99}O_{18}[0.33\text{ H}_2\text{O},0.08\text{ CO}_2]$$

prior to heating. After heating, through loss of the volatile channel occupants (and perhaps some Na),...
it became

\[ \text{K}_{0.01}\text{Na}_{0.03}(\text{Mg}_{8.91}\text{Fe}_{0.09}\text{Mn}_{0.01})\text{Al}_{3.98}\text{Si}_{5.61}\text{O}_{18} \]

Samples (~50 mg) of this tempered cordierite were sealed in platinum tubes and held for 4 weeks at 600°C under H\(_2\)O pressures of 0.5, 2, 4, 5 and 6 kbar. Others were similarly treated but in a CO\(_2\) atmosphere (from decomposition of Ag\(_2\)C\(_2\)O\(_4\)). Annealing times of 4 weeks were needed to achieve optical homogeneity within the crystal grains of the sample. The samples were subsequently checked by IR spectroscopy between 5000 cm\(^{-1}\) and 400 cm\(^{-1}\) using the KBr-powder technique. H\(_2\)O and CO\(_2\) were thus easily detected. H\(_2\)O displays symmetric and asymmetric stretching modes between 3570 cm\(^{-1}\) and 3700 cm\(^{-1}\) (Farrell and Newnham, 1967; Goldman et al., 1977). Gaseous CO\(_2\) possesses an asymmetric stretching vibration at 2349 cm\(^{-1}\). Hence, absorption at or near 2354 cm\(^{-1}\) for beryl (Wood and Nassau, 1967) and for cordierite (Farrell and Newnham, 1967) was attributed by them to CO\(_2\) trapped in the channels. For the optical studies, single grains from each run were mounted on a goniometer head and oriented on the spindle stage. Refractive indices were measured by the double variation (\(\lambda, T\)) method and \(2V\) was calculated from extinction data at 400, 666 and 900 nm applying the computer program EXCALIBR. A precise description of the optical methods used in this paper is given elsewhere (Selkregg and Bloss, 1980; Bloss, 1981). The precision of the refractive indices is within 0.0005 and \(2V\) within 0.5°. For the same crystal grains, cell dimensions were measured by the back-reflection Weissenberg method. For at least 60 indexed reflections (CuKa\(_1\), CuKa\(_2\) and CuK\(\beta\)), \(2\theta\) values were submitted to the least square program of Burnham (1962, 1965), as revised by L. Finger, which corrects for film shrinkage and absorption. Water in the samples was quantitatively determined by the equation of Medenbach et al. (1980) which relates change in average refractive index (upon heating to complete dehydration) to water content for Mg-cordierite. CO\(_2\) released from the sample at 1300°C C was determined using the same equipment and method as described in detail by Johannes and Schreyer (1980). CO\(_2\) and H\(_2\)O analyses lead to standard deviations below 5% of the total gas analyzed. The orientation of CO\(_2\) in natural cordierites was determined by polarized IR spectroscopy from oriented single crystal slabs kindly provided by D. S. Goldman (and described by Goldman et al., 1977). For these single crystal slabs, we compared the relative intensity of the asymmetric stretching absorption of CO\(_2\) at (ca.) 2350 cm\(^{-1}\) along the \(a\), the \(b\), and the \(c\) axis.

**Results**

The White Well cordierite, degassed by heating between 1200° and 1300°C in a slightly reducing atmosphere, developed oriented cracks but no color change or accompanying hematite reflections in its X-ray photographs. The cracks subsequently vanished for the H\(_2\)O-treated crystals, probably because H\(_2\)O enhances diffusion processes. By contrast, the CO\(_2\) treatment caused only partial healing of the cracks. IR spectra confirmed H\(_2\)O as the only volatile channel occupant in the H\(_2\)O-treated cordierites and CO\(_2\) as the only one in the CO\(_2\)-treated ones. The H\(_2\)O-containing cordierites show a strong absorption peak at 3690 cm\(^{-1}\) and weaker ones at 3575 cm\(^{-1}\) and 3630 cm\(^{-1}\). The CO\(_2\)-containing cordierites exhibit strong peaks at about 2350 cm\(^{-1}\). There is no significant difference in the frequencies of the absorption bands for channel CO\(_2\) and H\(_2\)O observed by us from those observed by Farrell and Newnham (1967) or Goldman et al. (1977). Absorption bands attributed to AlO\(_x\), SiO\(_x\), and MgO\(_x\) vibrations are consonant with the spectra observed by Langer and Schreyer (1969). Non-equilibrium conditions during hydration or carbonation were readily detected because the resultant crystals possessed higher refractive indices at their edges than at their cores. Moreover, such non-homogeneous crystals exhibited undulatory extinction because of the strong influence of H\(_2\)O and CO\(_2\) on the optic angle (Armbruster and Bloss, 1980). By contrast, crystals annealed at 600°C for 4 weeks displayed sharp extinction.

With increased gas in the channels, each refractive index increased significantly (Fig. 2, Table 1). As channel H\(_2\)O increased, \(\beta\) increased more rapidly than \(\gamma\) and much more so than \(\alpha\). Hence, \(2V_x\) decreases sharply. As channel CO\(_2\) increased, \(\gamma\) increased much more rapidly than \(\beta\) and \(\alpha\). Hence, \((\gamma - \alpha)\) and \(2V_x\) increased strongly with channel CO\(_2\) as first noted by Armbruster and Bloss (1980). CO\(_2\) and H\(_2\)O each affect the cell dimensions (Fig. 3, Table 2). Channel H\(_2\)O decreases \(a_0\) but it slightly increases \(b_0\) and \(c_0\) so that cell volume remains constant (within the range of error). As a result, the distortion index \(\Delta\), calculated by the simplified equation of Selkregg and Bloss (1980), namely

\[ \Delta = 1.094 (a_0 - b_0 \sqrt{3}) \]
The increase in refractive index of Mg-cordierite that results if \( \text{H}_2\text{O} \) occupies the channels (light lines, open rectangles) or if \( \text{CO}_2 \) occupies the channels (bold lines, solid rectangles). The edges of the rectangles correspond to two estimated standard deviations. The data for the channel evacuated cordierite are represented by concentric open and solid rectangles.

Avogadro number, and the electron polarizability for a molecular compound, where \( P \) is the polarizability averaged over all possible orientations of the electric vector. Although the equation is not strictly applicable to anisotropic crystals, Hartshorne and Stuart note that a principal refractive index (\( e, \omega, \alpha, \beta, \) or \( \gamma \)) may be substituted for \( n \). In such case \( P \) reasonably approximates the crystal's polarizability along the principal vibration direction associated with this principal index and may thus be symbolized \( P_e, P_\omega, P_\alpha, P_\beta \) or \( P_\gamma \), as the case may be. For example,

\[
P_e = \frac{e^2 - 1}{e^2 + 2} \cdot \frac{M}{\rho} \cdot \frac{3}{4\pi N}
\]

Actually, \( M/\rho \) represents the molar volume \( V_M \). For a crystal with known unit-cell volume \( (V_{uc}) \), if \( Z \) represents the formula units encompassed by \( V_{uc} \), then its molar volume can be calculated since

\[
V_M = \frac{N V_{uc}}{Z}
\]

Hence the preceding equation can be rewritten as, for example,

\[
P_e = \frac{3}{4\pi} \cdot \frac{e^2 - 1}{e^2 + 2} \cdot \frac{V_{uc}}{Z}
\]

For mineralogists this is convenient since \( V_{uc} \) and \( Z \) are more likely to be accurately known than density or molecular weight. The units for \( P \) will be those used for \( V_{uc} \) (usually \( \text{Å}^3 \)).

For the cordierites here studied, polarizabilities \( P_\alpha, P_\beta, \) and \( P_\gamma \) were calculated using their principal

| Table 1. Optical data of \( \text{H}_2\text{O} \) and \( \text{CO}_2 \) treated White Well cordierite at 600°C. |
|-----------------|-----------------|-----------------|-----------------|-----------------|-----------------|
| Pressure (kbar) | Gas (wt %) | \( n \) | \( \alpha \) | \( \beta \) | \( \gamma \) | \( (\gamma - \alpha) \) | 2\( \beta_\alpha \)* |
| ---             | ---         | 1.5235 | 1.5254 | 1.5275 | 1.529566 | 0.0040 | 87.0 |
| \( \text{H}_2\text{O} \) re-introduced into channels |
| 0.5            | 1.05       | 1.5303 | 1.5340 | 1.5362 | 1.53530 | 0.0059 | 64.6 |
| 2.0            | 1.68       | 1.5331 | 1.5389 | 1.5410 | 1.53833 | 0.0059 | 56.6 |
| 4.0            | 2.24       | 1.5386 | 1.5440 | 1.5450 | 1.54253 | 0.0064 | 50.0 |
| 6.0            | 2.56       | 1.5409 | 1.5464 | 1.5477 | 1.54503 | 0.0068 | 44.1 |
| \( \text{CO}_2 \) re-introduced into channels |
| 0.5            | 0.27       | 1.5238 | 1.5260 | 1.5287 | 1.52616 | 0.0049 | 86.8 |
| 2.0            | 0.91       | 1.5326 | 1.5305 | 1.5345 | 1.53043 | 0.0062 | 93.8 |
| 4.0            | 1.88       | 1.5351 | 1.5389 | 1.5410 | 1.53878 | 0.0133 | 105.0 |
| 5.0            | 2.61       | 1.5362 | 1.5410 | 1.5525 | 1.54323 | 0.0163 | 111.1 |
| 6.0            | 3.12       | 1.5388 | 1.5464 | 1.5582 | 1.54720 | 0.0187 | 112.4 |

* Average of measurements at 400, 666, and 900 nm.

decreased as \( \text{H}_2\text{O} \) content increased. With increased channel \( \text{CO}_2, a_0 \) first decreases but, beyond 2 weight percent \( \text{CO}_2 \), it increases; \( b_0 \) seemed unaffected by \( \text{CO}_2 \) but \( c_0 \) increased strongly. Consequently, \( \Delta \) decreased (up to 2 weight percent \( \text{CO}_2 \)), then increased slightly.

IR spectra for slabs cut from a crystal of natural cordierite confirmed the conclusions of Farmer (1974) that the linear \( \text{CO}_2 \) molecule aligns mainly in the (001) plane and, within this plane, chiefly along \( a \).

Hartshorne and Stuart (1970, p. 137) state the Lorentz-Lorenz equation for molecular refractivity \( R_M \) to be

\[
R_M = \frac{n^2 - 1}{n^2 + 2} \cdot \frac{M}{\rho} = \frac{4}{3} \pi NP
\]

where \( n, M, \rho, N \) and \( P \) respectively represent refractive index, molecular weight, density, the
Fig. 3. Variation in unit cell edges (A, B and C), distortion index Δ (D), and unit cell volume (E) as amount of channel occupant increases. The rectangles have the same significance as in Figure 2.

Table 2. Cell dimensions, volume and distortion index of H₂O and CO₂ treated White Well cordierite at 600°C.

<table>
<thead>
<tr>
<th>Pressure (kbar)</th>
<th>Gas (wt %)</th>
<th>a₀ [Å]</th>
<th>b₀ [Å]</th>
<th>c₀ [Å]</th>
<th>V [Å³]</th>
<th>Δ</th>
</tr>
</thead>
<tbody>
<tr>
<td>Channel-evacuated</td>
<td>---</td>
<td>---</td>
<td>17.076 (2)</td>
<td>9.722 (1)</td>
<td>9.346 (1)</td>
<td>1550.3 (3)</td>
</tr>
<tr>
<td>H₂O re-introduced into channels</td>
<td>0.5</td>
<td>1.05</td>
<td>17.106</td>
<td>9.725</td>
<td>9.343</td>
<td>1550.6</td>
</tr>
<tr>
<td></td>
<td>2.0</td>
<td>1.68</td>
<td>17.061</td>
<td>9.723</td>
<td>9.346</td>
<td>1550.4</td>
</tr>
<tr>
<td></td>
<td>4.0</td>
<td>2.24</td>
<td>17.060</td>
<td>9.725</td>
<td>9.346</td>
<td>1550.6</td>
</tr>
<tr>
<td></td>
<td>6.0</td>
<td>2.56</td>
<td>17.052</td>
<td>9.723</td>
<td>9.346</td>
<td>1550.1</td>
</tr>
<tr>
<td>CO₂ re-introduced into channels</td>
<td>0.5</td>
<td>0.27</td>
<td>17.069</td>
<td>9.721</td>
<td>9.345</td>
<td>1550.8</td>
</tr>
<tr>
<td></td>
<td>2.0</td>
<td>0.91</td>
<td>17.069</td>
<td>9.721</td>
<td>9.345</td>
<td>1550.8</td>
</tr>
<tr>
<td></td>
<td>4.0</td>
<td>1.88</td>
<td>17.059</td>
<td>9.722</td>
<td>9.349</td>
<td>1550.5</td>
</tr>
<tr>
<td></td>
<td>5.0</td>
<td>2.61</td>
<td>17.059</td>
<td>9.722</td>
<td>9.350</td>
<td>1551.1</td>
</tr>
<tr>
<td></td>
<td>6.0</td>
<td>3.13</td>
<td>17.062</td>
<td>9.722</td>
<td>9.356</td>
<td>1551.8</td>
</tr>
</tbody>
</table>

* Numbers in parentheses in the table represent the estimated standard deviation in terms of the least decimal place cited for the value to the left and for all other values in the column below. The precisions achieved in determining the gas contents are discussed in the text.
Table 3. Polarizabilities for the H$_2$O and CO$_2$ treated White Well cordierite

<table>
<thead>
<tr>
<th>Gas (wt %)</th>
<th>$P_\alpha$ (Å$^3$)</th>
<th>$P_\beta$ (Å$^3$)</th>
<th>$P_\gamma$ (Å$^3$)</th>
<th>$P_\delta$ (Å$^3$)</th>
</tr>
</thead>
<tbody>
<tr>
<td>Channel-evacuated</td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>1.05</td>
<td>28.601</td>
<td>28.767</td>
<td>28.866</td>
<td>28.745</td>
</tr>
<tr>
<td>1.68</td>
<td>28.813</td>
<td>28.983</td>
<td>29.077</td>
<td>28.958</td>
</tr>
<tr>
<td>2.24</td>
<td>29.973</td>
<td>29.215</td>
<td>29.259</td>
<td>29.149</td>
</tr>
<tr>
<td>2.36</td>
<td>29.067</td>
<td>29.317</td>
<td>29.370</td>
<td>29.251</td>
</tr>
<tr>
<td>CO$_2$ re-introduced into channels</td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>0.27</td>
<td>28.310</td>
<td>28.410</td>
<td>28.532</td>
<td>28.420</td>
</tr>
<tr>
<td>0.91</td>
<td>28.424</td>
<td>29.613</td>
<td>29.793</td>
<td>28.610</td>
</tr>
<tr>
<td>1.88</td>
<td>28.702</td>
<td>29.949</td>
<td>29.298</td>
<td>28.984</td>
</tr>
<tr>
<td>2.61</td>
<td>29.875</td>
<td>29.090</td>
<td>29.602</td>
<td>29.190</td>
</tr>
<tr>
<td>3.13</td>
<td>29.905</td>
<td>29.264</td>
<td>29.868</td>
<td>29.380</td>
</tr>
</tbody>
</table>

The increase in refractive indices as channel H$_2$O is reintroduced into channel-evacuated White Well cordierite (Table 3) is consonant with type I water. For this orientation, vectors drawn from the hydrogens to the nearest oxygens of the Al$_2$Si$_4$O$_{18}$ rings would represent directions of maximum (but weak) polarization of these oxygens by the hydrogens. Such directions, being more nearly parallel than perpendicular to (001), would hence favor, as Table 3 shows, a greater rate of increase of the indices $\beta$ and $\gamma$ than of $\alpha$ with increased channel water (since $X = c$). The long distance (3.4Å) between the water oxygen and the nearest oxygen of the Al$_2$Si$_4$O$_{18}$ ring precludes the existence of hydrogen bonding because, when hydrogen bonds exist, the associated O-O distances usually equal 2.7–2.8Å (Vinogradov and Linell, 1971). Moreover, Langer and Schreyer (1976) point out that the slight energy decrease of about 80 cm$^{-1}$ (0.2 kcal) for the IR-stretching modes of H$_2$O in cordierite compares to those for H$_2$O vapor, this again indicating unlikelihood of strong hydrogen bonds, these latter usually being in the range of 5 kcal.

The linear CO$_2$ molecule (4.96Å) has been observed with its elongation mostly normal to the c axis (Farrell and Newnham, 1967) and only to a smaller degree parallel to c (Farmer, 1974). Our IR results confirm this. Such observations of the orientation of CO$_2$ (and H$_2$O) in natural specimens may be complicated because specific CO$_2$ and H$_2$O orientations may result from interaction with other channel occupants (or each other). On the other hand, the small shifts of the CO$_2$ bands in IR spectra for cordierite as compared to those of CO$_2$ in the gaseous state suggest only minor interaction with the silicate framework. The marked increase in the refractive index $\gamma$, as channel CO$_2$ is reintroduced into a channel-evacuated White Well crystal (Table 3), suggests that CO$_2$ is dominantly aligned parallel to $a (= Z)$. The refractive index $\alpha$ displays the least

Fig. 4. Variation of the Lorentz-Lorenz polarizability $P$ (in Å$^3$) for White Well cordierite crystals as their channels are increasingly occupied by H$_2$O (hollow squares) and by CO$_2$ (solid squares). Estimated standard deviations are not implied by the size of the squares.
rate of increase with channel CO₂. This suggests that CO₂ is dominantly aligned perpendicular to c (since \(X = c\)), even when not parallel to \(a\), and thus agrees with the IR results.

Elongation of the CO₂ molecule parallel to \(a\) (type \(a\) orientation) avoids strong interaction with the oxygens that form the channel walls and with any other CO₂ molecules in the cavities above or below (Fig. 5A). The most probable position for CO₂, if oriented parallel to \(c\) (type \(c\) orientation), is with its carbon centered within the Al₂Si₄O₁₈ ring (Fig. 5B) so that its oxygen atoms extend into the cages. In such case, however, only every second cavity could be occupied by CO₂ because \(c/c\) equals 4.7Å whereas the length of CO₂ is 4.96Å. Type \(c\) orientation might be favored by a low degree of channel filling, perhaps coupled with presence of cations (Na\(^+\), etc.) in adjacent available sites.

With increased re-introduction of channel CO₂, \(a_0\) initially decreased but then increased, \(b_0\) remained constant, and \(c_0\) increased significantly (Table 2, Fig. 3). With increased channel H₂O, \(a_0\) decreased, \(b_0\) perhaps increased slightly, and \(c_0\) initially increased but then levelled off. Channel H₂O and channel CO₂ both decreased the distortion index \(\Delta\) (Table 2).

Under the same \(P-T\) conditions, fewer CO₂ molecules than H₂O molecules enter cordierite’s channels. At 600°C and 6 kbar, approximately twice as many H₂O molecules as CO₂ molecules enter (so that the number of oxygens entering is about the same in each case). Presumably, the elongate CO₂ molecule enters the channel like a needle entering a tube of slightly larger diameter. After reaching a cavity, the CO₂ molecule must rotate 90° to become parallel to \(a\), its favored orientation. With heating to expel CO₂ from the channels, any CO₂ molecule aligned along \(a\) must rotate 90° in order to diffuse outward along the channels or else decompose to smaller breakdown products.

Puzzlingly, at 600°C and 5 kbar, we observed 2.61 wt.% CO₂ to enter the White Well cordierite whereas Johannes and Schreyer (1980) observed that, under these same \(P-T\) conditions only 1 wt.% CO₂ entered finely powdered synthetic Mg-cordierite. In each case the same equipment and methods were used to determine CO₂. Possible synthetic cordierites possess a greater frequency of channel offsets or dislocations (by twin, domain, or grain boundaries). Also, Johannes and Schreyer (1980), using two different types of synthetic Mg-cordierites, noted a pronounced effect of the starting material on the amount of CO₂ (but not H₂O) that entered. For natural cordierites, on the other hand, the effect of channel Na or of Fe content upon entry of CO₂ is not as yet known. A catalytic interaction between Na and CO₂ might favor incorporation of CO₂ into the channels. These conjectures aside, the internal consistency of our optical data, which yield Lorentz–Lorenz polarizabilities that plot quite linearly relative to wt.% CO₂ (Fig. 4), corroborate a value of 2.61 wt.% CO₂ for the grain held at 600°C, 5 kbar in a CO₂ atmosphere.

Knowledge of the effect of CO₂ and H₂O on the
optical properties of cordierites may enable us (Armbruster and Bloss, in preparation) to estimate the absolute and/or relative amounts of H₂O and CO₂ in natural cordierites. Ultimately, it is hoped, such data may provide estimates of the composition of the fluid phase with which a cordierite was last in equilibrium.

Acknowledgments

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References

Tang, T., and Ghose, S. (1972) Nuclear magnetic resonance of 1H and 27Al and Al-Si order in low cordierite Mg₆Al₄Si₁₄O₃₈·n H₂O. Journal of Chemical Physics, 56, 3329-3332.

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