

The crystal structure of mercury(II)phosphate, $\text{Hg}_3(\text{PO}_4)_2$

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Auszug

Die Struktur des Quecksilber(II)phosphates, $\text{Hg}_3(\text{PO}_4)_2$, wurde mittels symbolischer Addition und Fourierrmethoden bestimmt. Durch Verfeinerung aus 1491 unabhängigen Zählrohrintensitäten wurde ein R -Faktor von 0,043 erhalten. Die Kristalle sind monoklin, die Raumgruppe ist $P2_1/c$, mit vier Formeleinheiten pro Elementarzelle und den folgenden Zellparametern: $a = 9,737(2) \text{ \AA}$, $b = 11,466(2) \text{ \AA}$, $c = 6,406(1) \text{ \AA}$ und $\beta = 99,51(2)^\circ$.

Die Quecksilberatome sind auf beinahe lineare Weise zweifach-koordiniert, die Abstände Quecksilber-Sauerstoff schwanken zwischen $2,06(1) \text{ \AA}$ und $2,13(1) \text{ \AA}$ und die O—Hg—O-Winkel zwischen $163,4(4)^\circ$ und $169,9(5)^\circ$. Alle Sauerstoffatome gehören zu Phosphattetraedern. Jedes Quecksilberatom ist an je ein Sauerstoffatom von zwei Phosphattetraedern gebunden und jede Phosphatgruppe besitzt Bindungen zu drei Quecksilberatomen.

Die Struktur ist aus unendlichen, einander durchdringenden und gefalteten Netzen mit der Formel $[\text{Hg}_3(\text{PO}_4)_2]_n$ aufgebaut. Die Netze werden aus endlosen

—Hg—O—P—O—Hg—Ketten, miteinander durch zusätzliche Quecksilberatome verbunden, gebildet. Die Hg—O-Abstände zwischen den Netzen sind $\geq 2,42(1) \text{ \AA}$.

Abstract

The structure of mercury(II)phosphate, $\text{Hg}_3(\text{PO}_4)_2$, has been determined by symbolic addition and Fourier methods and refined to an R value of 0.043 on the basis of 1491 independent counter intensities. The crystals are monoclinic, space group $P2_1/c$, with four formula units in a unit cell of the dimensions $a = 9.737(2)$, $b = 11.466(2)$, $c = 6.406(1) \text{ \AA}$ and $\beta = 99.51(2)^\circ$.

The mercury atoms are two-coordinated in a nearly linear way; the mercury to oxygen distances vary between $2.06(1)$ and $2.13(1) \text{ \AA}$, and the O—Hg—O angles between $163.4(4)$ and $169.9(5)^\circ$. All oxygen atoms belong to phosphate tetrahedra. Each mercury atom is bonded to one oxygen atom of each of two phosphate tetrahedra and each phosphate group is bonded to three mercury atoms.

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The structure is built up of infinite, interpenetrating puckered nets of formula $[\text{Hg}_3(\text{PO}_4)_2]_n$, formed from endless $-\text{Hg}-\text{O}-\text{P}-\text{O}-\text{Hg}-$ chains, fused by additional mercury atoms. The $\text{Hg}-\text{O}$ distances between the nets are $\geq 2.42(1)$ Å.

Introduction

In connection with several studies on mercury salts containing pyramidal or tetrahedral anions (BJÖRNLUND, 1971, 1974; AURIVILLIUS, 1972), it was found of interest to investigate some mercury phosphates. This paper deals with the structure of $\text{Hg}_3(\text{PO}_4)_2$. Work is in progress on a compound of formula $\text{Hg}_2(\text{H}_2\text{PO}_4)_2$, containing monovalent mercury.

Experimental

Single crystals of $\text{Hg}_3(\text{PO}_4)_2$ were prepared according to two different methods. First a synthesis suggested by KLEMENT and HASELBECK (1964) was tried, the procedure somewhat modified by the present authors, using a more dilute solution of mercury(II)oxide in phosphoric acid (3 g HgO in 50 ml concentrated H_3PO_4). Another method is to use an acidified solution of mercury(II)acetate and phosphoric acid (10 ml 0.3M $\text{Hg}(\text{C}_2\text{H}_3\text{O}_2)_2$, 10 ml 2.5M HNO_3 , 10 ml conc. H_3PO_4). The crystallization proceeds very slowly in the last case, but gives a more well-formed product. The crystals thus obtained were in the form of thin, colourless plates. Weissenberg photographs showed the crystals to be monoclinic. The systematically absent reflections were $h0l$ with $l \neq 2n$ and $0k0$ with $k \neq 2n$, the extinctions consistent with the space group $P2_1/c$. X-ray powder photographs were taken in a Guinier-Hägg focusing camera with $\text{CuK}\alpha_1$ radiation, using aluminium as internal standard ($a = 4.04934$ Å). The preliminary values of the cell dimensions were refined by least-squares calculations. The density was determined from the loss of weight in benzene. Some crystal data are given in Table 1.

Table 1. *Crystal data*

$\text{Hg}_3(\text{PO}_4)_2$, Formula weight 791.8
Monoclinic, $P2_1/c$
$a = 9.737(2)$, $b = 11.466(2)$, $c = 6.406(1)$ Å,
$\beta = 99.51(2)^\circ$,
$V = 705.4$ Å ³ , $Z = 4$.
$D_m = 7.32$, $D_x = 7.45$ g · cm ⁻³
$\mu(\text{MoK}\alpha) = 669$ cm ⁻¹

A single crystal with the dimensions $0.20 \times 0.07 \times 0.007$ mm was used for the intensity data collection on a Pailred linear diffractometer, using $\text{MoK}\alpha$ radiation, monochromatized by reflection off the (002) planes of a graphite crystal, the monochromator angle being 6.08° . The crystal used was mounted along its longest edge, which coincides with the direction of the crystallographic b axis. The reflections of the layer lines $h0l$ — $h11l$ were collected for the copper range, $(\sin \theta)/\lambda \leq 0.65$, using the equi-inclination and ω -scan technique with a scan rate of $1.0^\circ/\text{min}$. The scan range was 4.0° for all reflections. The stationary background counts were measured for 40 sec at each end of the scan interval. The aperture size of the detector was 2.0° . As a check of the electronic stability during the period of data collection, the intensity of two standard reflections, 400 and 050, were measured at regular intervals. The fluctuation of their intensities was random, and the values of $(I_{\max} - I_{\min})/I_{\max}$ were less than 6%. A total of 2232 independent reflections were recorded ($h \geq 0$, $k \geq 0$). 209 reflections for which the two measured background values differed more than 3.09 times the estimated standard deviations of their difference were omitted. The integrated peak counts I were calculated from the total integrated peak counts, the background counts and the counting time in the usual way. 532 reflections with $I \leq 2.58\sigma(I)$ were considered unobserved. The values of $\sigma(I)$ were based on counting statistics. The remaining 1491 reflections were corrected for Lorentz, polarization and absorption effects. The transmission factors, evaluated by numerical integration, varied from 0.04 to 0.61.

Structure determination and refinement

The positions of the heavy atoms were determined by symbolic addition methods. The asymmetric part of the E map showed three maxima of nearly the same height, indicating the positions of the atoms $\text{Hg}(1)$ — $\text{Hg}(3)$ of the unit cell. Least-squares refinement followed by difference Fourier syntheses revealed the positions of the phosphorus and oxygen atoms. The positional parameters, the isotropic temperature factors and the inter-layer scale factors were then included in a full-matrix least-squares refinement, minimizing $\sum w_i(|F_o| - |F_c|)^2$ with weights $w_i = [\sigma^2(F_o) + a|F_o|^2]^{-1}$. A suitable value for a was found from an analysis of the weighting scheme. The discrepancy factors R and R_w , defined by $R = \sum ||F_o| - |F_c|| / \sum |F_o|$ and $R_w = [\sum w_i(|F_o| - |F_c|)^2 / \sum w_i |F_o|^2]^{1/2}$, converged to $R = 0.064$ and $R_w =$

$= 0.083$. When anisotropic temperature factors were introduced for all atoms in a new refinement, the R value fell to 0.043 and R_w to 0.051. In the last cycle of refinement the value of a was 0.001. At the end of the refinement the value of S was 0.99. The expression for S (goodness of fit) is $S = [\sum w_i(|F_o| - |F_c|)^2/(m - n)]^{1/2}$, where

Table 2. *Atomic coordinates obtained in the final least-squares refinement*
Estimated standard deviations are given in parentheses

Atom	x	y	z
Hg(1)	1.04137(7)	0.62072(5)	0.67502(8)
Hg(2)	0.69987(7)	0.39641(6)	0.82674(8)
Hg(3)	0.65409(7)	0.67005(6)	0.41226(9)
P(1)	0.6109(4)	0.4010(4)	0.3051(5)
P(2)	0.8972(4)	0.8666(3)	0.5987(5)
O(11)	0.551(1)	0.525(1)	0.249(2)
O(12)	0.716(1)	0.407(1)	0.511(2)
O(13)	0.698(1)	0.356(1)	0.139(2)
O(14)	0.492(1)	0.316(1)	0.317(2)
O(21)	1.034(1)	0.798(1)	0.583(2)
O(22)	0.783(1)	0.778(1)	0.632(2)
O(23)	0.916(1)	0.940(1)	0.804(2)
O(24)	0.860(1)	0.941(1)	0.399(2)

Table 3. *Thermal parameters β_{ij} with estimated standard deviations in parentheses.*
The expression used is $\exp [-(\beta_{11} \cdot h^2 + \beta_{22} \cdot k^2 + \beta_{33} \cdot l^2 + 2\beta_{12} \cdot hk + 2\beta_{13} \cdot hl + 2\beta_{23} \cdot kl)]$. The β_{ij} values are multiplied by 10^5 for Hg, 10^4 for P and 10^3 for O. The root-mean-square components, R_i , of thermal vibration along principal axes of the ellipsoids of vibration are also given

Atom	β_{11}	β_{22}	β_{33}	β_{12}	β_{13}	β_{23}	R_1	R_2	R_3
Hg(1)	276(6)	110(5)	734(12)	23(4)	30(6)	33(4)	0.126 Å	0.083 Å	0.113 Å
Hg(2)	288(6)	282(5)	458(11)	29(4)	49(6)	-22(5)	0.139	0.096	0.114
Hg(3)	341(6)	228(6)	794(12)	114(4)	171(7)	121(5)	0.155	0.095	0.115
P(1)	19(4)	12(3)	38(7)	-4(2)	5(4)	-2(3)	0.104	0.077	0.088
P(2)	14(3)	6(3)	48(7)	0(2)	2(4)	-1(3)	0.100	0.062	0.081
O(11)	3(1)	3(1)	13(3)	0(1)	-1(2)	1(1)	0.183	0.108	0.145
O(12)	1(1)	4(1)	6(2)	0(1)	1(1)	1(1)	0.162	0.047	0.112
O(13)	3(1)	3(1)	5(2)	0(1)	2(1)	0(1)	0.142	0.088	0.124
O(14)	3(1)	5(1)	8(2)	-2(1)	4(1)	-1(1)	0.199	0.026	0.132
O(21)	2(1)	1(1)	12(3)	1(1)	1(1)	0(1)	0.155	0.075	0.119
O(22)	1(1)	3(1)	8(2)	-1(1)	1(1)	-1(1)	0.158	0.046	0.119
O(23)	2(1)	2(1)	8(2)	0(1)	2(1)	-1(1)	0.138	0.090	0.109
O(24)	1(1)	2(1)	8(2)	0(1)	-1(1)	0(1)	0.136	0.044	0.108

Table 4. *Selected interatomic distances and angles in the structure of $\text{Hg}_3(\text{PO}_4)_2$. Estimated standard deviations are given in parentheses. For notations of the atoms, cf. Table 2*

Mercury to oxygen distances		O—Hg—O angles	
within the nets		O(21)—Hg(1)—O(23)	163.4(4)°
Hg(1)—O(21)	2.11(1) Å	O(13)—Hg(2)—O(12)	169.9(5)
Hg(1)—O(23)	2.11(1)	O(11)—Hg(3)—O(22)	163.9(5)
Hg(2)—O(13)	2.06(1)		
Hg(2)—O(12)	2.06(1)		
Hg(3)—O(11)	2.13(1)		
Hg(3)—O(22)	2.12(1)		
between the nets			
Hg—O	$\geq 2.42(1)$ Å		
Distances and angles within the PO_4 tetrahedra			
P(1)—O(11)	1.55(1) Å	P(2)—O(21)	1.56(1) Å
—O(12)	1.53(1)	—O(22)	1.55(1)
—O(13)	1.55(1)	—O(23)	1.55(1)
—O(14)	1.53(1)	—O(24)	1.53(1)
Mean values	1.541(6)		1.545(6)
O(1 <i>i</i>)—O(1 <i>j</i>)	2.43(2)— 2.57(2) Å	O(2 <i>i</i>)—O(2 <i>j</i>)	2.43(2)— 2.58(2) Å
O(11)—P(1)—O(12)	109.2(7)°	O(21)—P(2)—O(22)	108.3(6)°
O(11)—P(1)—O(13)	111.8(7)	O(21)—P(2)—O(23)	110.2(6)
O(11)—P(1)—O(14)	109.5(8)	O(21)—P(2)—O(24)	108.3(6)
O(12)—P(1)—O(13)	104.0(6)	O(22)—P(2)—O(23)	103.3(6)
O(12)—P(1)—O(14)	113.1(6)	O(22)—P(2)—O(24)	113.9(6)
O(13)—P(1)—O(14)	109.3(7)	O(23)—P(2)—O(24)	112.7(6)

m denotes the number of observations and n the number of parameters varied. Thus the errors seem to be correctly estimated. All parameter shifts were less than 5% of the estimated standard deviations. Corrections for extinction and for anomalous dispersion were not included in the calculations, however. A final three-dimensional difference synthesis was calculated, with the contributions of all atoms subtracted. The highest remaining peak had an approximate height of $3\text{e}/\text{\AA}^3$. The scattering factors were those of CROMER and WABER

[illegible]

k	$ P_0 / P_c $	k	$ P_0 / P_c $	k	$ P_0 / P_c $	k	$ P_0 / P_c $	k	$ P_0 / P_c $	k	$ P_0 / P_c $	k	$ P_0 / P_c $	k	$ P_0 / P_c $	k	$ P_0 / P_c $	k	$ P_0 / P_c $
4	173	170	20	295	285	5	73	73											
5	232	232				6	90	90	1										
6	93	96				7	57	66											
7	92	85	1	63	50														
8	136	138	3	283	261	0	k	2											
10	63	70	4	55	46	0	239	242	9	107	113	4	63	61					
11	81	82	5	185	177	1	152	151	9	105	105	5	84	89	1	k	4		
9	k	0	6	109	103	2	138	141	10	74	73	6	204	205					
0	49	51	7	149	143	3	338	311	11	75	68								
1	80	83	9	185	179	6	195	205											
2	76	90	10	140	135	7	478	498	1	114	130	4	k	3					
4	101	107	5	k	1	8	205	207	3	246	236	4	k	3					
6	161	164				9	48	50	3	129	144	1	85	84	6	194	194		
7	19	51	5	562	522	10	130	125	4	240	243	7	253	243	5	96	100		
8	134	136	6	273	269	11	308	325	8	148	141	3	185	184	9	87	84	7	
9	112	118	9	74	84	2	k	11	96	189	5	46	43	9	122	121			
10	115	115	11	132	127	1	k	2	6	39	30	10	153	156	11	k	4		
11	50	51				0	517	549							0	236	223	8	k
10	k	0	1	246	251	2	188	188	1	39	31	9	72	72					
0	59	63	4	159	156	3	104	103	4	167	153								
2	95	90	5	395	378	4	375	393	6	115	108	5	k	3	2	115	115	0	k
6	187	189	6	111	118	5	104	107	10	69	73	1	110	106	3	55	58	2	62
8	44	43	9	111	6	52	52					2	181	176	4	53	51	3	82
10	77	81	9	327	312	10	k	2	3	373	376	3	5	11	115	106	9	k	5
11	71	70	10	192	181	8	44	45	0	137	135	5	373	376	5	213	218	5	102
11	k	0	7	k	1	9	119	123	1	76	73	7	65	67	3	51	36	5	124
						10	69	70	4	92	84	7	63	63	8	63	61	8	116
1	185	198	1	189	173	2	k	1	2	7	61	61	8	89	82	9	72	76	9
3	145	152	2	181	162	3	124	127	0	149	152	10	110	104	9	270	269	10	125
5	98	101	3	267	243	11	k	2	10	123	123	3	k	4	1	162	164	10	k
5	98	1																	

(1965) for neutral mercury and of HANSON *et al.* (1964) for neutral phosphorus and oxygen. The final positional and thermal parameters, together with the root-mean-square components, are given in Tables 2

and 3 and selected interatomic distances and angles in Table 4. Observed and calculated structure amplitudes are presented in Table 5.

Description and discussion of the structure

As seen from Table 4 the mercury atoms are each coordinated to two oxygen atoms, the Hg—O distances varying from 2.06(1) to 2.13(1) Å and the O—Hg—O angles from 163.4(4) to 169.9(5)°. The values of the distances are in good agreement with those found in the infinite —O—Hg—O— chains of orthorhombic HgO [2.04(3), 2.07(3) Å] (AURIVILLIUS, 1956, 1964), and in the —Hg—O—Cr—O—Hg— chains of $\text{HgCrO}_4 \cdot \frac{1}{2} \text{H}_2\text{O}$ [2.055(1), 2.064(1) Å] (AURIVILLIUS and STÅLHANDSKE). The values of the O—Hg—O angles are more distorted in the present compound, however [179.5(1.1)° in HgO; 179.95(5)° in $\text{HgCrO}_4 \cdot \frac{1}{2} \text{H}_2\text{O}$]. A dominant feature of the structure of $\text{Hg}_3(\text{PO}_4)_2$ is thus the frequently occurring two-covalency of mercury(II).

The structure is built up of infinite puckered nets of formula $[\text{Hg}_3(\text{PO}_4)_2]_n$. Two such nets, related by $\bar{1}$, run through the unit cell. The nets are interpenetrating and the Hg—O distances between them are $\geq 2.42(1)$ Å.

The nets are formed from fused endless chains —Hg—O—P—O—Hg— in the following way: All oxygen atoms of the structure belong to phosphate tetrahedra; one of them consists of the atoms P(1), O(11)—O(14), and the other of the atoms P(2), O(21)—O(24). Each mercury atom is bonded to one oxygen atom of each of two phosphate groups, while three oxygen atoms of each phosphate group are bonded to mercury. The atoms Hg(1), bridging between the atoms O(21) and

O(23), form separate endless —Hg—O—P—O—Hg— chains related to each other through $\bar{1}$. They run approximately parallel to [010]. In the same way, infinite chains, nearly parallel to [001], are formed as the atoms Hg(2) bridge the atoms O(12) and O(13) (*cf.* Table 2 for labelling of the atoms). The endless puckered nets are formed from —O(21)—Hg(1)—O(23)— and —O(12)—Hg(2)—O(13)— chains, fused by additional mercury atoms, Hg(3), which in turn are bonded to the atoms O(22) and O(11). The fourth corner of each tetrahedron, the atom O(24) or O(14), is not bonded to any mercury atom at short distance. The Hg—O distances within the nets are 2.06(1)—2.13(1) Å.

If the shortest mercury to oxygen distance between the nets of 2.42(1) Å is also considered as a coordination distance, the structure

may be described as built up of a three-dimensional network of formula $[\text{Hg}_3(\text{PO}_4)_2]_n$. A stereoscopic view of the structure is given in Fig. 1.

The phosphate tetrahedra are rather regular (Table 3), the mean $\text{P}(1)-\text{O}$ and $\text{P}(2)-\text{O}$ distances being 1.541(6) and 1.545(6) Å, respectively. These values are in good agreement with mean $\text{P}-\text{O}$ values calculated from data given for other phosphates, *e.g.* Li_3PO_4 [1.546(3) Å] (KEFFER *et al.*, 1967), $\text{Mg}_3(\text{PO}_4)_2$ [1.527(3) Å] (NORD and

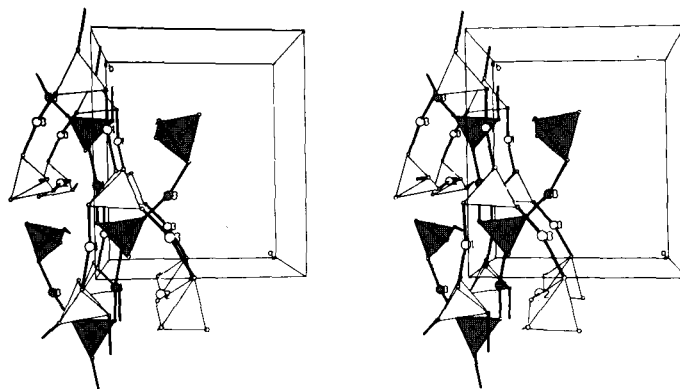


Fig. 1. Stereoscopic view along the c axis, showing two endless puckered nets of formula $[\text{Hg}_3(\text{PO}_4)_2]_n$, related by $\bar{1}$. The three mercury atoms are indicated as 1, 2 and 3. The nets are shadowed and unshadowed in the drawing. The c axis points against the reader

KIERKEGAARD, 1968) and $\text{Zn}_3(\text{PO}_4)_2$ [1.533 Å] (CALVO, 1965). There is, however, a small tendency towards elongation of the $\text{P}-\text{O}$ bonds for the $\text{P}-\text{O}-\text{Hg}$ bridging oxygen atoms, the largest difference in the bonds being about 3σ in the distances (0.036 Å). Much larger deformations of the tetrahedral groups are found in *e.g.* $\text{HgCrO}_4 \cdot \frac{1}{2} \text{H}_2\text{O}$, where the elongation of the $\text{Cr}-\text{O}$ bonds for the two $\text{Cr}-\text{O}-\text{Hg}$ bridging

oxygen atoms in the infinite $-\text{Hg}-\text{O}-\text{Cr}-\text{O}-\text{Hg}-$ chains amount to 0.12 Å ($\approx 60\sigma$ in the distances). In $\text{Hg}(\text{OH})_2 \cdot 2 \text{HgSO}_4 \cdot \text{H}_2\text{O}$, one oxygen atom of each sulfate tetrahedron bridges between sulfur and mercury in the limited chains $\text{O}(\text{SO}_4)-\text{Hg}-\text{O}(\text{OH})-\text{Hg}-\text{O}(\text{OH})-\text{Hg}-\text{O}(\text{SO}_4)$, giving an elongation of 0.08 Å ($\approx 8\sigma$ in the distances).

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