

Far infrared and Raman spectroscopic investigations of lunar materials from Apollo 11, 12, 14, and 15

CLIVE H. PERRY, D. K. AGRAWAL, E. ANASTASSAKIS,
R. P. LOWNDES, and N. E. TORNBORG

Solid State Spectroscopy Laboratory, Physics Department,
Northeastern University, Boston, Massachusetts 02115

Abstract—We have studied the elastic and inelastic light scattering of twelve lunar surface rocks and eleven lunar soil samples from Apollo 11, 12, 14, and 15, over the range 20–2000 cm^{-1} . The phonons occurring in this frequency region have been associated with the different chemical constituents and are used to determine the mineralogical abundances by comparison with the spectra of a wide variety of terrestrial minerals and rocks. Kramers-Kronig analyses of the infrared reflectance spectra provided the dielectric dispersion (ϵ' and ϵ'') and the optical constants (n and k). The dielectric constants at $\sim 10^{11}$ Hz have been obtained for each sample and are compared with the values reported in the 10^2 – 10^6 Hz range. The emissivity peak at the Christianson frequencies for all the lunar samples lie within the range 1195–1250 cm^{-1} ; such values are characteristic of terrestrial basalts. The Raman light scattering spectra provided investigation of small individual grains or inclusions and gave unambiguous interpretation of some of the characteristic mineralogical components.

INTRODUCTION

WE HAVE STUDIED the interaction of electromagnetic radiation with lunar materials using elastic and inelastic light scattering techniques. These investigations have many important applications as they identify the vibrational modes associated with the different chemical constituents and can be used to determine mineralogical abundances and give structural information. The data can be used for compositional mapping of the lunar surface and the uniqueness of the method suggests the possibility of performing remote infrared petrology using spectral matching techniques (Lyon, 1963; Aronson *et al.*, 1967; Perry and Lowndes, 1970). The establishment of the relative chemical constituents and their relation to the different rock and soil types on the lunar surface is made by comparison with spectra of a wide variety of terrestrial minerals and rocks.

The infrared studies were accomplished from measurements (mainly conducted at 300°K) of the specular component of the reflectivity over the thermal frequency range 20–2000 cm^{-1} (500–5 μ wavelength). The data were subjected to a Kramers-Kronig analysis as described previously (Perry *et al.*, 1971) and the dispersion of the real and imaginary parts of the dielectric function (ϵ' , ϵ''), the refractive index n , and the absorption coefficient α were obtained. These fundamental optical constants are independent of the method of measurement and can be used directly to simulate emittance spectra of lunar materials for comparison with future remote sensed infrared spectra and radiometric measurements of the lunar surface (Mendell and Low, 1971).

The far-infrared data yield values of the dielectric constant and loss tangent at $\sim 10^{11}$ Hz. It is of interest to extrapolate these values of the dielectric function to

lower frequencies to provide useful information for electrical prospecting of subsurface layers in the 10^6 – 10^9 Hz range (Simmons *et al.*, 1971; Gold *et al.*, 1972; Howard and Tyler, 1972). Comparison can then be made with the values obtained in the 10^2 – 10^6 Hz range using standard capacitance techniques provided that these measurements have been made on dry material (Collett and Katsuba, 1971; Chung *et al.*, 1971; Chung, 1971; Chung and Westphal, 1972; Lowndes *et al.*, 1972).

The inelastic (Raman) light scattering techniques for determining the frequency, structural, and chemical properties of the constituents of lunar samples offers a complementary tool. In contrast to the infrared measurements which provide macroscopic information over large areas, this technique, by virtue of the small size ($< 50 \mu$ diameter) of the focused laser beam, allows the investigation of small individual grains or inclusions and has provided unambiguous interpretation of some of the characteristic mineralogical components (Perry *et al.*, 1971; White *et al.*, 1971; Estep *et al.*, 1972; Fabel *et al.*, 1972).

EXPERIMENTAL

Infrared measurements

We have obtained the infrared specular reflectance spectra of the following bulk samples which were in the form of polished butt ends: 10058,56, 12002,186, 12008,23, 12009,48, 12065,115, 12073,42, 14301,20, 14307,18, 14310,184, 14313,51, 14321,98, and 15426,2. The samples were polished to obtain one smooth flat face having an area of at least 0.25 cm^2 .

Pressed disks $\sim 0.1 \text{ cm}$ thick of the fines samples were examined in a similar manner. The samples were prepared in a standard 0.5 cm diameter die with pressures of about $5 \times 10^9 \text{ dynes/cm}^2$ being applied for approximately 5 minutes. The densities of all the samples ranged from 2.4 – 2.6 g/cm^3 . The following fines samples were examined: 14141,37, 14161,36, 14163,31, 14230,95, 14259,38, 15021,159, 15071,60, 15091,65, 15221,71, 15471,68, and 15531,63.

Dry nitrogen flushed grating spectrometers (PE-301, PE-180, and PE-521) with specular reflectance attachments (angle of incidence 10° – 15°) were used for the studies over the frequency range 200 – 4000 cm^{-1} (2.5 – 50μ). From 20 – 250 cm^{-1} a Fourier transform Michelson type interferometer with a liquid-helium-cooled germanium detector was used (Perry *et al.*, 1966).

Raman measurements

Identifiable Raman spectra of a wide variety of individual grains and glassy inclusions were obtained on a number of lunar samples. Various excitation radiations were used, namely Ar^+ 4880 and 5145 \AA and He/Ne 6328 \AA , and some resonance effects were noted. The Raman scattered light was analyzed using a Spex 1401 double spectrometer and detected using photon counting techniques. Both oblique incidence and back-scattering sample geometries were employed. The power absorbed by the samples often represented a restriction on the Raman technique as the opaque regions showed a definite tendency to decompose under the influence of the focused laser beam ($< 50 \mu$ diameter with a power density of $5 \times 10^4 \text{ W/cm}^2$). Both Stokes and anti-Stokes components of the spectrum were examined at room temperature.

Generally, the Raman spectra were of lesser complexity than their infrared counterparts and comparison with terrestrial mineral spectra provided unambiguous interpretation. The effects due to the lowering of site symmetry in the $(\text{SiO}_2)^{x-}$ group (Griffith, 1969), variation of Fs content in the pyroxenes, and Fa content in the olivines were demonstrated by frequency shifts in the spectra (Perry *et al.*, 1971; Estep *et al.*, 1971; 1972).

Dielectric studies

The low-frequency dielectric constants were determined through three terminal capacitance measurements in the frequency range 10^2 – 10^5 Hz. The capacitance measurements were recorded on a

General Radio 1615A transformer ratio arm bridge used in conjunction with a General Radio 1232A tuned amplifier and null detector. The capacitances were determined to an accuracy of $\pm 0.1\%$ at the lower frequencies and $\pm 0.4\%$ at 10^5 Hz. The variable temperature cell used to measure the dielectric constants has been described previously (Lowndes and Martin, 1969).

RESULTS AND DISCUSSION

The infrared reflectance spectra from $10\text{--}1700\text{ cm}^{-1}$ of some representative lunar rock samples from Apollo 11, 12, and 14 are shown in Fig. 1. From these spectra the real and imaginary parts of the dielectric constant were calculated and are plotted in Fig. 2. In Fig. 3 the infrared reflectance spectra of three Fra Mauro breccias 14301,20; 14307,18 and 14313,51, and a typical Apollo 14 soil 14161,36 (bulk sample) are displayed. The spectra of these breccias can be compared with those of an Apollo 12 mare breccia (12073,42) and an Apollo 15 breccia (15426,2) from the Apennine Front

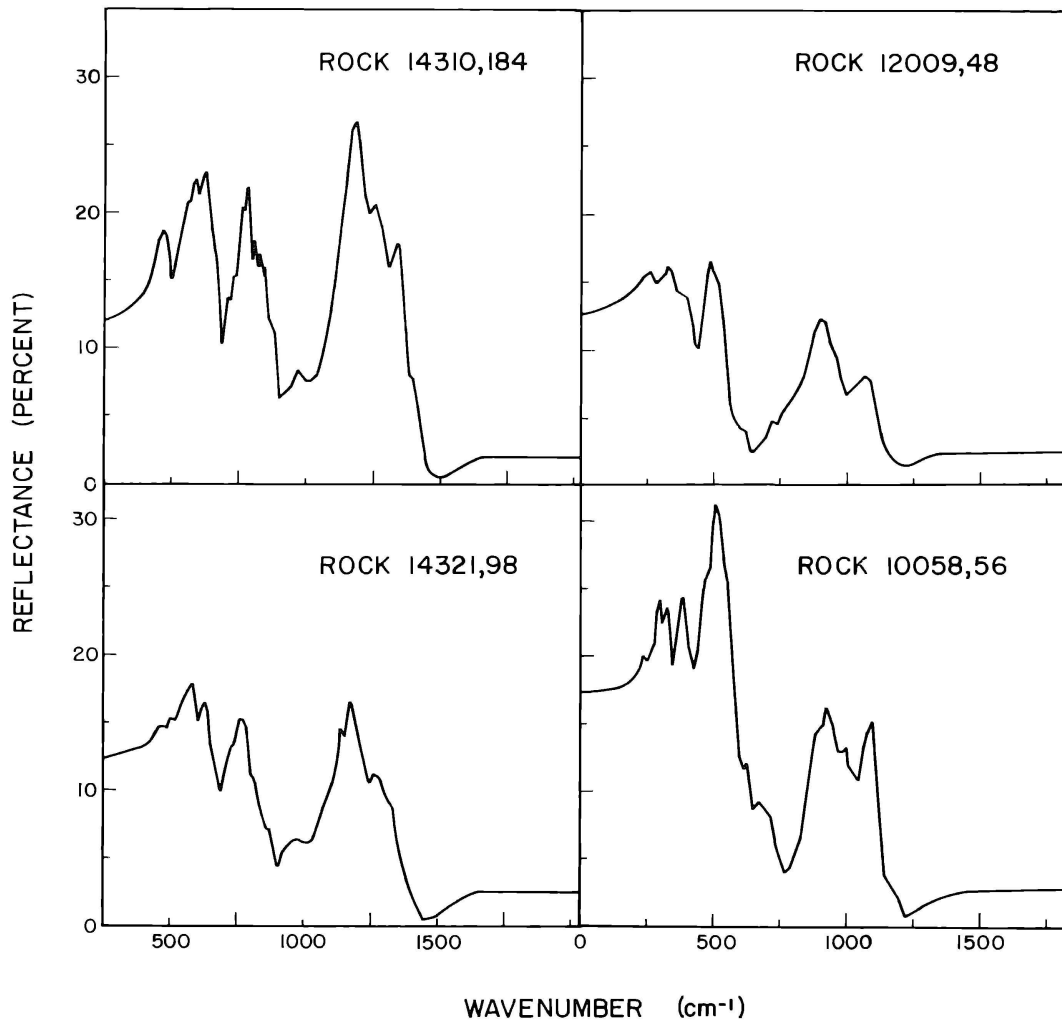


Fig. 1. The infrared reflectance spectra of rock 14310,184, 12009,48, 14321,98, and 10058,56, from 0 to 1700 cm^{-1} .

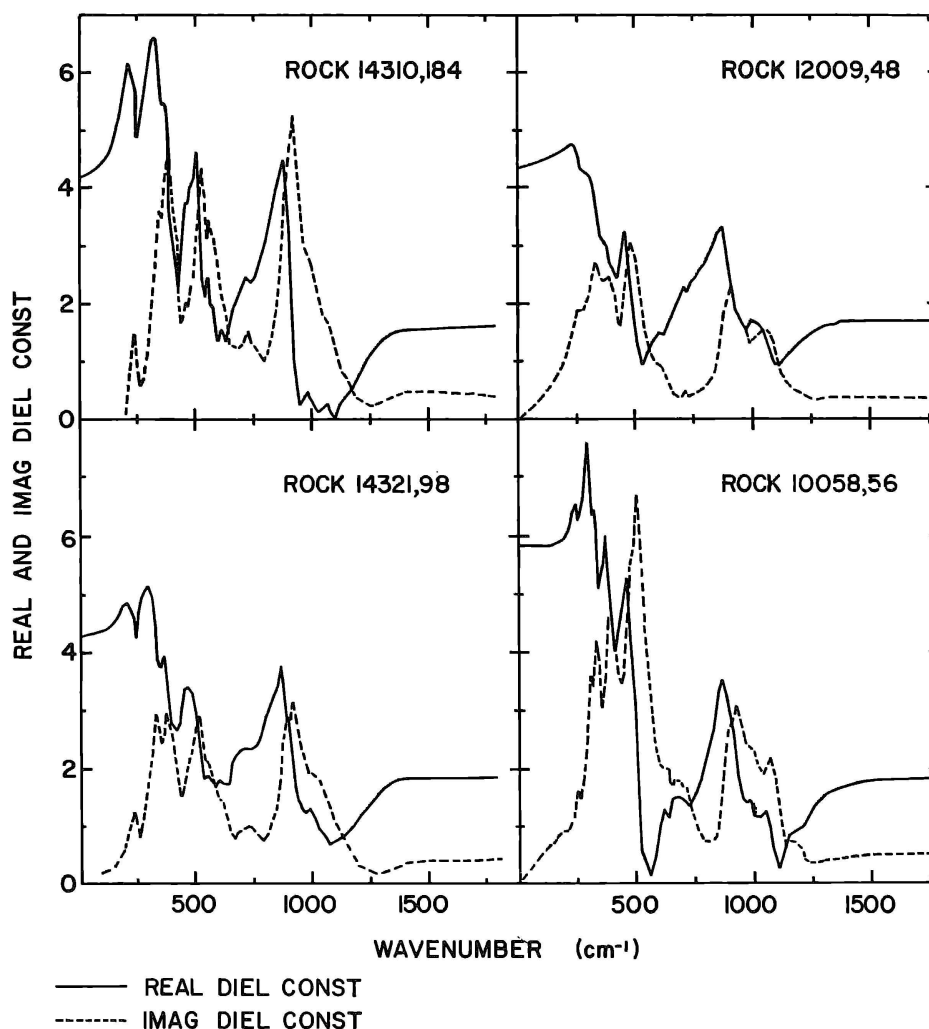


Fig. 2. The real and imaginary parts of the dielectric constant for rock 14310,184, 12009,48, 14321,98, and 10058,56, from 0 to 1700 cm^{-1} .

shown in Fig. 4. Figure 5 comprises the spectra of five Apollo 14 fines (14163,31; 14161,36; 14259,38; 14230,95 (core sample); and 14141,37 (Cone Crater)) and one powdered rock sample 14321,263 (Cone Crater). The reflectance spectra of six Apollo 15 fines are shown in Fig. 6. The samples investigated were 15091,65 and 15221,71 from Station 2 at the foot of the Apennine Front, 15071,60 from Station 1 (Elbow Crater ejecta); 15021,159 from the LM area, 15471,68 from Dune Crater (Station 4), and 15531,63 from Station 9a at the rim of Hadley Rille.

The spectra shown in the first six figures represent the bulk mineralogical and chemical properties of the materials averaged over the area of the sample examined (25–100 mm^2). They form a basis for the comparison of these samples with each other and with known terrestrial materials. In Figs. 7 and 8 are shown some examples of the results of the Raman investigations. Because of differing selection rules in the scattering process, this technique is complementary to that of infrared spectroscopy in the sense that it provides information unobtainable by the latter. Furthermore, be-

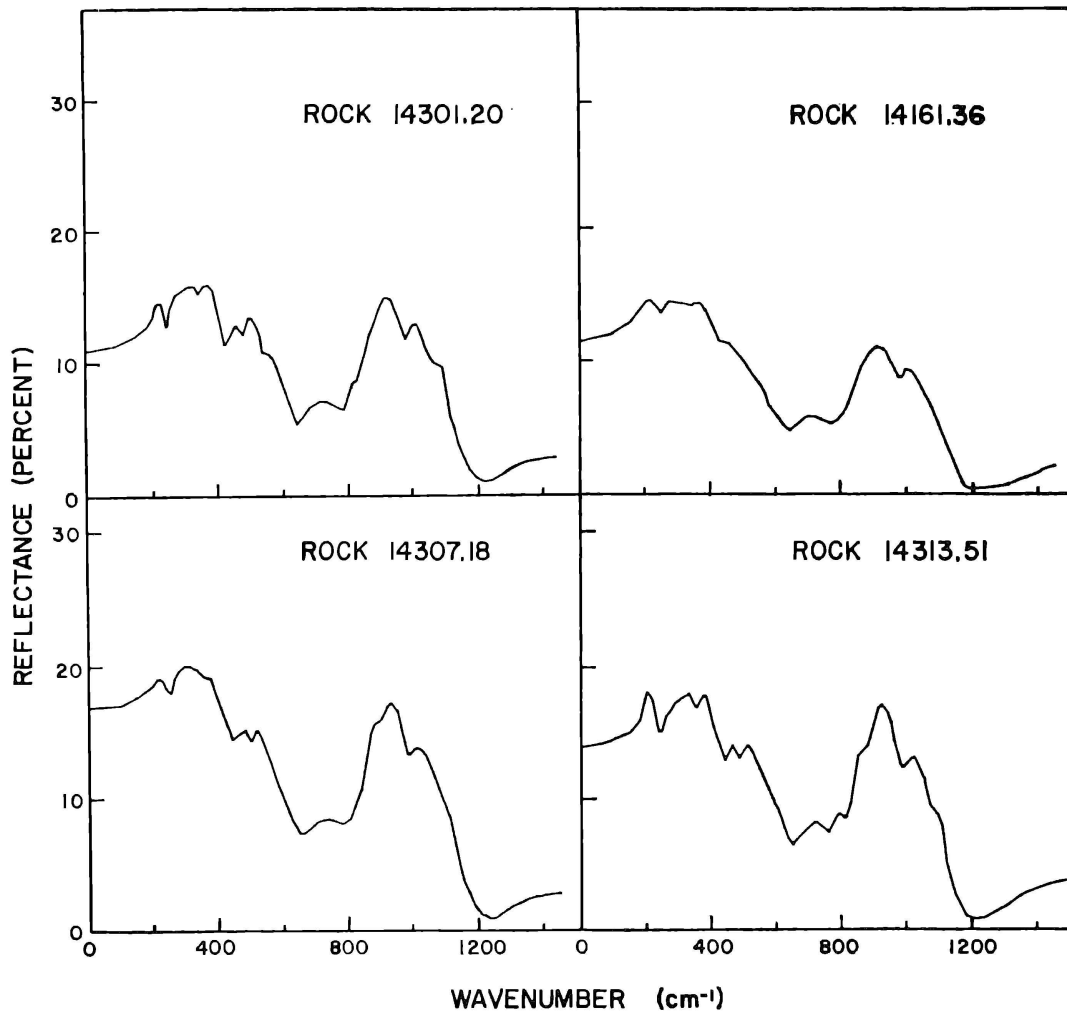


Fig. 3. The infrared reflectance spectra from 0 to 1400 cm^{-1} of three Fra Mauro breccias 14301,20, 14307,18, and 14313,51, and a typical Apollo 14 soil 14161,36 (bulk sample).

cause of the small size of the sample area investigated ($\sim 2 \times 10^{-5} \text{ cm}^2$), individual grains and inclusions may be studied, and the resulting spectra are usually simpler and the mineralogical identification is more straightforward. In Fig. 7, for example, the Raman spectra show some of the variations of lunar pyroxenes in samples 10058,56, 12002,186, 12065,115, and 14310,76 (Perry *et al.*, 1971, and as discussed by Estep *et al.*, 1972). The systematic frequency shift of the stretching vibration at $\sim 1015 \text{ cm}^{-1}$ (e.g., Estep *et al.*, 1972) and the bending vibration at $\sim 650 \text{ cm}^{-1}$ in these samples can be correlated with the ferrosilite content. In 14310,76 (with $\text{Fs}_{25}\text{-Fs}_{30}$) the latter band is almost a resolved doublet indicating the presence of both ortho- and clinopyroxene. Sample 10058,56 contains mostly clinopyroxene ranging from pigeonite to augite ($\text{Fs}_{20}\text{-Fs}_{40}$). Samples 12002,186 and 12065,115 show broader, weaker spectra and the evaluation is less well defined but both again probably contain clino- and orthopyroxene. The Raman spectrum of a typical green glass sphere in sample 15426,2

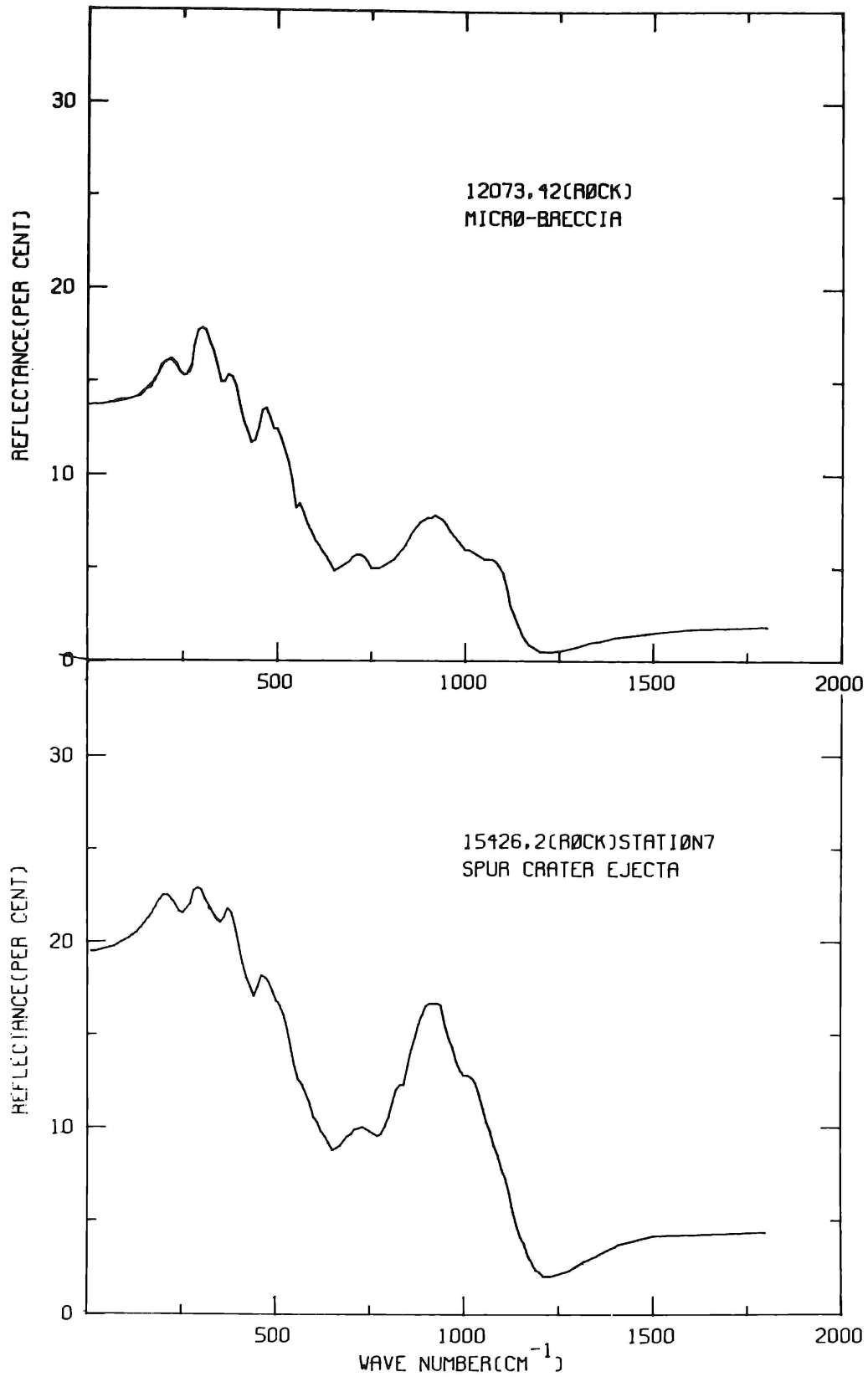


Fig. 4. The infrared reflectance spectra of an Apollo 12 mare breccia 12073,42 and an Apollo 15 breccia 15426,2 from the Apennine Front.

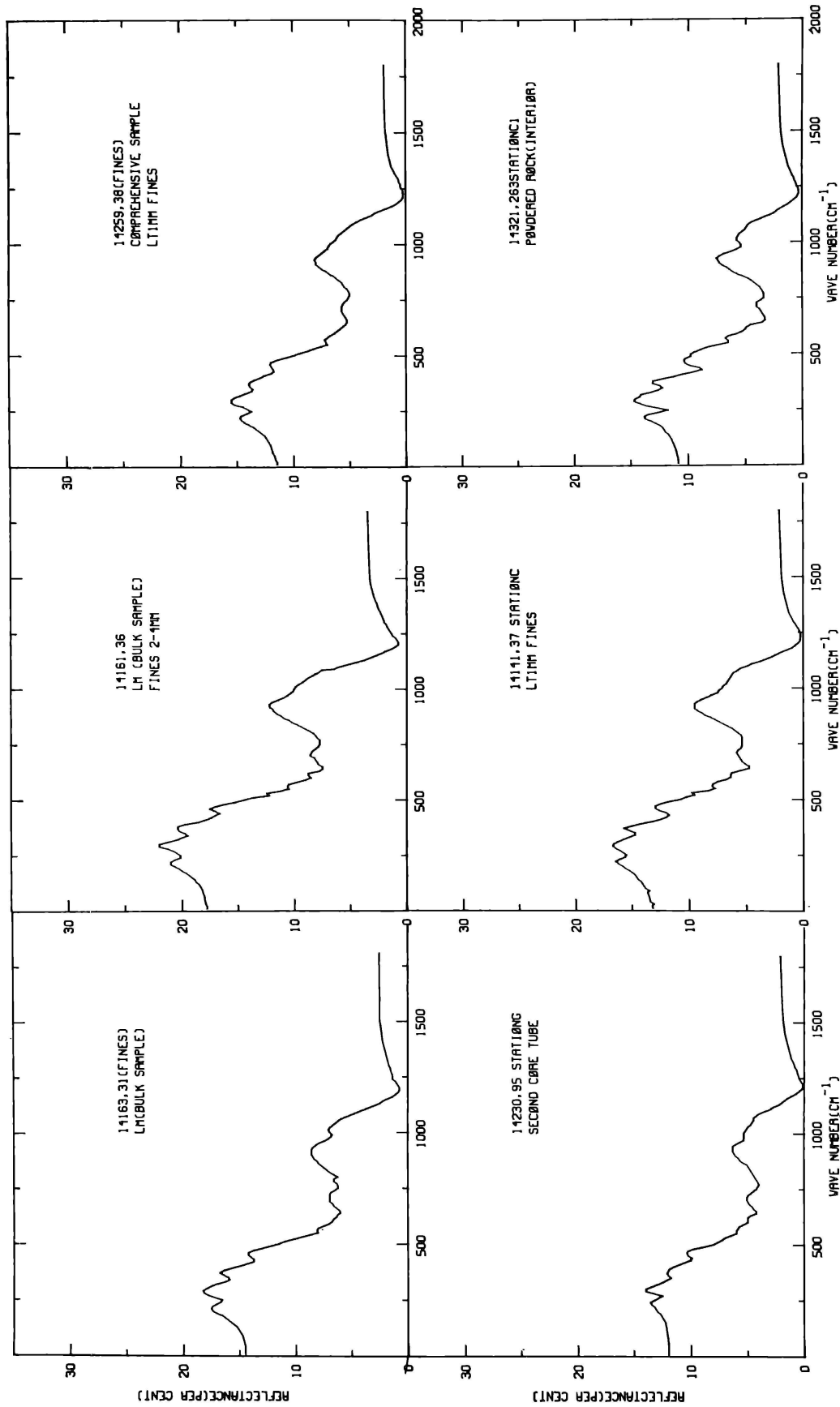


Fig. 5. The infrared reflectance spectra of five Apollo 14 fines (14163,31, 14161,36, 14259,38, 14230,95 (core sample), and 14141,37 (Cone Crater) and one powdered rock sample 14321,263 (Cone Crater) from 0 to 1700 cm⁻¹.

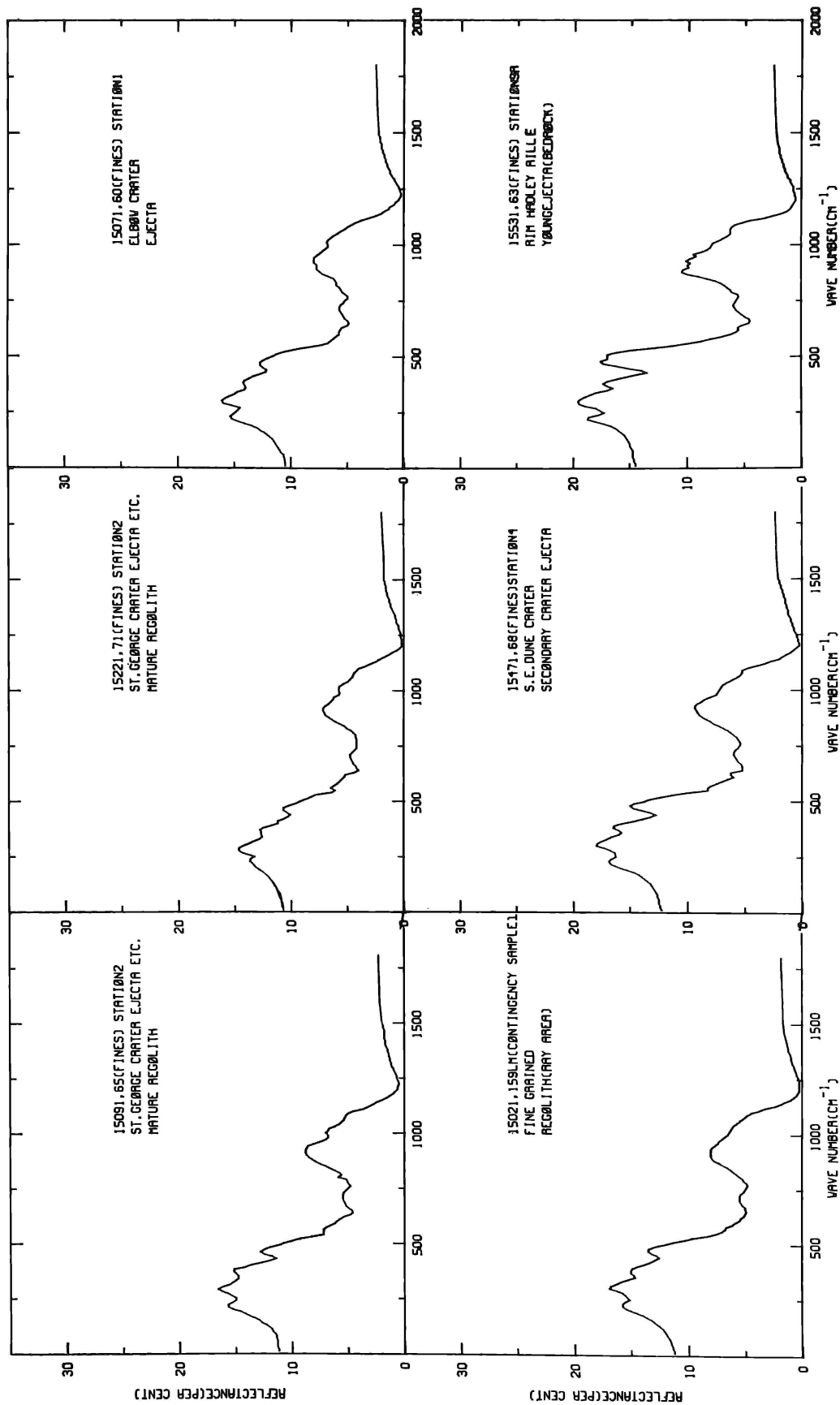


Fig. 6. The infrared reflectance spectra of six Apollo 15 fines (15091,65, 15221,71, 15071,60, 15021,159, 15471,68, and 15531,63) from 0 to 1700 cm^{-1} .

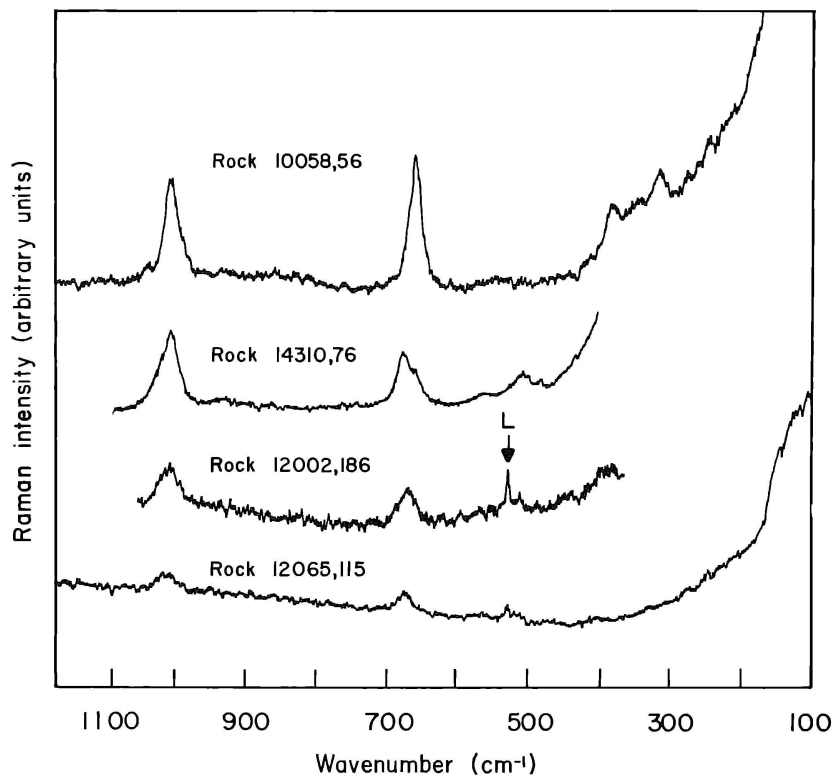


Fig. 7. The Raman spectra of lunar pyroxene inclusions in samples 10058,56, 14310,76, 12002,186, and 12065,115 showing the variation in the frequencies of bands that can be related with the ferrosilite content. The L refers to an extraneous laser line.

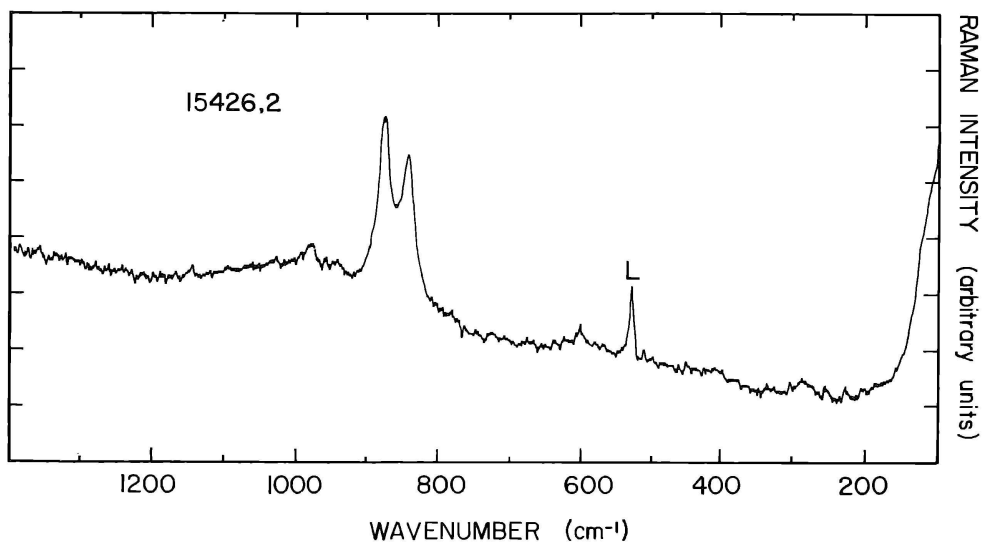


Fig. 8. The Raman spectrum of a typical green glass sphere in sample 15426,2. The strong doublet in the region of 850 cm^{-1} shows the presence of a particle that has undergone partial devitrification to olivine. The L refers to an extraneous laser line.

Table 1. A summary of the lunar rock samples investigated from the Apollo 11, 12, 14, and 15 missions. ω_c is the Christianson frequency corresponding to the emissivity peak and ϵ_0 is the dielectric constant at far infrared frequencies ($\sim 10^{11}$ Hz). The major chemical abundances are taken from the references cited.

Sample Number	Rock Classification	Lattice Modes	M-O-M bend oxygen vibs.	M-O stretchings	Combinations	ω_c (cm ⁻¹)	ϵ_0 at 10 ¹¹ Hz	Major chemical abundance (%)					§	
								SiO ₂	Al ₂ O ₃	MgO	FeO	CaO		TiO ₂
14310,184	Basaltic cryst. rock					1250	5.2	50	20	8	8	11	1	§
14321,98	Breccia (cone crater)					1205	4.3	49	16	12	11	8	1	
12002,186	Basalt					1220	3.9							
14301,20	Coherent Breccia					1230	4.0	48	16	10	12	10	2	*
14307,18	Coherent Breccia					1230	5.5	47	16.5	10	12	11	2	**
14313,51	Coherent Breccia					1215	4.8							
15426,2	Breccia					1220	6.5							
12073,42	Microbreccia					1240	2.7	41	15	11	17	11.5	3	†
12008,23	Cumulate					1200	6.7	Major phase, ilmenite						
12009,48	Prophyritic Basalt					1240	4.3	41	11	12.5	20	10	3	†
10058,56	Ophitic Basalt					1220	5,8	41	11	6	17	12	11	*
12065,115	Variolitic Basalt					1200	3.8	39	12	9	22	12	4	†

WAVE NUMBERS (cm⁻¹)

x Charles *et al.* (1971). † LSPET (1970). § LSPET (1972). * Hubbard and Gast (1972). ** Agrell *et al.* (1972).

is shown in Fig. 8. The strong doublet in the region of 850 cm^{-1} indicates devitrification to olivine (with $\text{Fa}_{20}\text{-Fa}_{40}$) of a contaminating particle either on or imbedded in the sphere. Raman spectra of lunar glasses have also been reported by White *et al.* (1971) and Fabel *et al.* (1972). However, Fabel *et al.* (1972) report a large number of closely spaced bands at low frequencies in some of their samples. We have observed similar bands in some of our samples but substitution of different laser excitation radiation revealed them to be grating ghosts.

Tables 1, 2, and 3 summarize all the lunar rock and soil samples that we have investigated from the Apollo 11, 12, 14, and 15 missions. A pictorial comparison is shown in Table 1 between the frequencies and oscillator strengths of the stretching, bending, and lattice modes for each sample. For more detailed information reference should be made to Perry *et al.* (1971, Figs. 3 and 4), and Tables 2 and 3.

Several researchers (Conel, 1969; Salisbury *et al.*, 1970) have suggested that the frequency of the peak in the emissivity can be used as a diagnostic feature for determining the silica abundance in fine-powder emissivity spectra. This peak is directly related through the dielectric function ($\epsilon' = n^2 - k^2$ and $\epsilon'' = 2nk$, where n is the real part of the refractive index and k is the extinction coefficient) to the high-frequency minimum in the reflectivity and in fact our reflectance spectra are simply inverted emission spectra referenced to blackbody emissivity. For powder spectra both curves are complicated by particle size and packing fraction because of additional diffuse scattering mechanisms (Aronson, 1972). This frequency (known as the Christianson frequency and referred to as ω_c in Table 1) occurs in the region where $n \approx 1$ and $k \approx 0$. Perry *et al.* (1971) have shown that for lunar materials ω_c varies over the narrow range from about 1195 cm^{-1} ($\sim 40\%$ SiO_2) to about 1260 cm^{-1} ($\sim 50\%$ SiO_2). These results for terrestrial materials correspond to the intermediate to ultrabasic range as reported by Conel (1969) and Salisbury *et al.* (1970). The emissivity peak for 14259 soil by Logan *et al.* (1972) at 1220 cm^{-1} is in good agreement with the ω_c value

Table 2. Comparison of infrared transverse optical phonon frequencies for lunar rock samples from Apollo 14 and 15.

Suggested assignment Combination	14310,184 1065(w)	14321,198	14301,20	14307,18	14313,51	15426,2 1040(w)
M-O stretch					1015(m)	
ν_3''	990(m)	1010(w) 980(w)	1010(w)	1005(m)		
ν_3'	920(s)	920(m)	940(m)	930(s)	925(s)	
ν_3		890(m)				850(m)
ν_1						810(w)
Cation-oxygen vibrations	725(w)	620(w)	745(w)	745(w)	730(w)	710(w) 640(w) 610(w) 570(w)
M-O-M bend	530(m)		530(m)			
ν_4		520(m)		525(m)	520(m)	
ν_4'	470(m)		475(m)	485(m)	470(s)	450(w)
ν' , ν_4''''	390(m)	390(m)	390(m)	380(m)	380(m)	370(w)
ν_2''	345(m)	335(m)	345(w)		345(m)	340(w)
Lattice modes	240(w)	240(w)	240(w)		225(w)	280(s)

Table 3. Comparison of the infrared optical phonon frequencies, ω_c (Christianson frequency) and ϵ_0 (dielectric constant for sample densities 2.4–2.6 g/cm³) for Apollo 14 and 15 lunar fines.

	Fines (LM)	2-4mm Fines (LM)	Comp. < 1 mm fines	Core sample	Station C < 1 mm fines	Station C1 Powdered rock
Suggested assignment Combinations	14163,31	14161,36	14259,38	14230,95 1060	14141,37 1050	14321,263
M–O stretch	1020 860	1020 930 810	1010 930 870	1010 900	900	1020 920
Cation-oxygen lattice vibrations	780 690	690 610	700	780 660 610	660 600	800 690
M–O–M bend	560 450 370	560 540 450 360	570 540 440 370	560 460 370	550 530 510 430 360	560 490 450 360 330
Lattice modes	290	290 230	300 200	280 230	260 220	270 220
ω_c (cm ⁻¹)	1195	1210	1225	1215	1230	1225
ϵ_0	4.9	5.9	4.0	4.2	4.5	3.8
	Station 2 St. George Crater ejecta	Station 1 Elbow Crater ejecta	Ray area (LM)	Station 4 Dune Crater	Station 9a Hadley Rille	
Suggested assignment Combinations	15091,65 1090	15221,71 1100	15071,60	15021,159 1040	15471,68 1100	15531,63 1070
M–O stretch	1000 880	1000 890	1000 920 870 820	890	1000 950 860	960 930 910 870
Cation-oxygen lattice vibrations	790 670 600 550	790 670 610 550	780 670 610 580	660	770 680 620 560	790 700 630
M–O–M bend	440 360	440 370	460 370	470 370	450 370	510 450 380
Lattice modes	280 220	280 230	290 230	300 210	280 230	280 220
ω_c (cm ⁻¹)	1220	1220	1225	1220	1210	1205
ϵ_0	4.0	4.2	3.8	4.0	4.2	4.9

(1225 cm^{-1}) obtained from our reflectance measurements as given in Table 3. Birkebak and Dawson (1972) obtained a value of 1210 cm^{-1} for 14163 soil having a density of 1.6 g/cm^3 which is within experimental error of our values of 1205 cm^{-1} for 14163,31 (LM sample) with a density of $\sim 2.5 \text{ g/cm}^3$. It is interesting to note that the infrared emissivity measurement of surface features in various lunar regions from balloon observations by Salisbury *et al.* (1970) and Murcray *et al.* (1970) vary from 1195 to 1235 cm^{-1} again within this narrow region. The 14310,184 sample, which we estimate to have a high anorthite content, compares with noritic gabbros and anorthitic basalts in terms of the Christianson frequency ($\omega_c \approx 1250 \text{ cm}^{-1}$), but there are marked anomalies in the series, as can be seen in Table 1, where ω_c is compared with the major chemical abundances. However, portions of the specimens used to obtain the elemental analyses are not the same as the ones used to measure the ω_c values, and variations can be expected within the same parent sample. Consequently, correlation of the ω_c values with SiO_2 wt.% may be better than the data presented in Table 1, although particle size and packing of the material also effect the scattering components and can cause variations in the measured Christianson frequency.

In Tables 1 and 3 we have also shown the values of the far-infrared dielectric constant at $\sim 10^{11}$ Hz which were calculated from the Kramers-Kronig analysis. Our values for the lunar rocks appear to be in general agreement with Ade *et al.* (1971) far-infrared measurements at 337 μ where a cyanogen laser was used. For sample 12063,96 (type A basalt) and for 10065,30 (breccia) they obtain values of 6.25 and 3.7, respectively. Based on the results of Perry *et al.* (1971) one may estimate that both these samples contain 25–30 mole percent of titani-ferrous oxide. Sample 10017,64, another Apollo 11 type A basalt, has a high dielectric constant value of 7 which compares with our measured value of 6.7 for 12008,23, and consequently it probably has a comparable or slightly higher ilmenite content (i.e., 30–35%). Comparison can be made with the dielectric response determined by capacitance measurements at low frequencies (10^2 – 10^6 Hz) by Chung *et al.* (1971), Collett and Katsuba (1971), Chung and Westphal (1972), and Lowndes *et al.* (1972), and at 430 MHz by Gold *et al.* (1972). The disagreement with previous values of Chung *et al.* (1971) appears to be due to their samples not being dry (Chung *et al.*, 1972). Our value for a 14310 sample of 5.2 (10^{11} Hz) and 5.99 (10^2 – 10^6 Hz) compares favorably with Chung and Westphal's (1972) value of 6–9 and the Gold *et al.* (1972) value of 6.6. However, there are still striking differences in sample 12002 where we obtained a dielectric constant (at 10^{11} Hz) of 4 which is to be compared with Chung *et al.*'s (1971) result of 8–10 and Collett and Katsuba's (1971) value of 9 at low frequencies. Similar differences were observed in our own measurements in 12008,23 where the dielectric constant varies between 6.7 at 10^{11} Hz and 33 at 10^6 – 10^2 Hz. As both these samples contain free iron and other high-conductivity materials, they may give rise to additional relaxation mechanisms or free carrier absorption which could account for this strong dispersion of the dielectric constant.

In an earlier report, Perry *et al.* (1971) have found that sample 12008,23 has a strong vibration at $\sim 890 \text{ cm}^{-1}$. The high value of the measured static dielectric constant of this sample indicates that it has a very high ($> 30\%$) FeO content and it has also been reported that the major phase of the parent rock may be ilmenite (Warner,

1970). Sample 12009,48 has a lower FeO content ($\sim 20\%$) and the stretching vibration is now situated at $\sim 920\text{ cm}^{-1}$. With the increase in plagioclase content, as for example in 12002,186 ($\sim 30\%$ plagioclase), there is a continued shift of this band to higher frequencies ($\sim 930\text{ cm}^{-1}$), and in 12065,115 (breccia), with about 50-50 plagioclase-pyroxene content, the trend is further indicated ($\sim 960\text{ cm}^{-1}$).

In Fig. 1 we show some of the greatest variations that we have observed in the reflection spectra of lunar samples. The 14310,184 spectrum has features that clearly demonstrate that it is a high calcium oxide, alumina, and silica basalt and it has the highest ω_c frequency compared with the other lunar samples examined. It has high stretching vibrations at $\sim 990\text{ cm}^{-1}$ and $\sim 920\text{ cm}^{-1}$ and bending vibrations at $\sim 470\text{ cm}^{-1}$ and 530 cm^{-1} . The lattice band at 240 cm^{-1} is characteristic of a high plagioclase material and is indicative of its high anorthite ($\text{An}_{90}\text{-An}_{100}$) content. For comparison the Cone Crater sample (14321,98) from Station C1 (which has comparable silica, alumina, and calcium oxide content) shows a lowering of the frequency of the band at $\sim 500\text{ cm}^{-1}$. This band appears to be related to the slightly higher orthopyroxene content. The spectrum of 12009,48 with its reduced silica and alumina content shows a further lowering of the bending vibration to $\sim 490\text{ cm}^{-1}$ which appears associated with the increase in FeO content (see Table 1). Generally, as the dominance of the framework silicates change gradually to the more open chain silicates of the pyroxenes, a lowering of the internal frequencies may be expected and is apparently observed. Finally in Fig. 1, the new features in the spectrum of sample 10058,56, which has FeO, SiO_2 , Al_2O_3 , and CaO contents comparable with those of 12009,48, indicate that they can be directly associated with the increasing TiO_2 present and the subsequent reduction in MgO.

The three Apollo 14 breccias (or fragmental rocks) examined, namely, 14307,18, 14313,51, and 14301,20, show unmistakable spectral similarities, as seen in Fig. 3. All are typical of the Fra Mauro feldspar-rich breccias and their spectra, although slightly enhanced in intensity, possess features that correspond closely with those observed in the Apollo 14 fines sample of which the bulk sample 14161,36 from the LM vicinity is a typical example.

From Table 1 it is seen that these Fra Mauro breccias have lower silica and alumina contents than sample 14310 which is more likely of highland origin since it is rich in plagioclase (anorthosites). However, they are intermediate between the highland material and the mare material of Apollo 11 and 12 and these findings have been confirmed in the x-ray fluorescence experiment by Adler *et al.* (1972). These breccias are presumably representative of the ejecta blanket deposited by the impact that formed the Imbrium Basin and contain materials from the upper levels of the crust. The Apollo 14 soils also show considerable homogeneity in the spectra (see Fig. 5) and this is confirmed in the analyses (Compston *et al.*, 1972), where the average soil content is 48% SiO_2 , 17% Al_2O_3 , 9% MgO, 10% FeO, 11% CaO, and 2% TiO_2 . These major chemical abundances are very similar to those found in the Apollo 14 breccias (see Table 1) and the spectral similarities between Fig. 3 and Fig. 5 indicate that the Fra Mauro soils were mainly produced by local impacts and consist of particulate debris from the rocks with a small admixture of other material (see, for example, Adams and McCord, 1972). From the infrared and Raman frequencies it

would appear that these samples contain plagioclase An_{70} – An_{90} , olivine Fa_{20} – Fa_{35} , and pyroxene Fs_{25} – Fs_{35} , and are in general agreement with the report on 14313 (Floran *et al.*, 1972).

The Cone Crater breccia sample from Station C, 14321,98, has spectral similarities that lie between 14310 and the other three Apollo 14 rock samples. It has the highest magnesium content and lowest calcium content and although the silica abundance is similar to 14310 its alumina and FeO contents are more like the Fra Mauro breccias. It may in fact be part of an earlier regolith that suffered from the Imbrian ejecta and was later ejected during the Cone Crater event. The increase in iron content in 14321 over the 14310 samples is exhibited in a shifting to lower frequencies of the band $\sim 500\text{ cm}^{-1}$. Its lower Christianson frequency probably results from its high MgO content. This reverse behavior of ω_c is apparent in 12009,48 which again has a comparable abundance of MgO even though the SiO_2 content is less and the FeO content is increased.

In Figs. 5 and 6 we can see a direct comparison between the Apollo 14 and 15 soil samples. Samples 14259,38, which is representative of the regolith, and 14163,31, which is from a depth of a few centimeters, both gave very similar spectra. Sample 14161,36, which consists of 2–4 mm fines, showed some slight sharpening of the reflectivity peaks. The sample 14230,95 from the second core tube again showed general overall spectral similarities and it was concluded that this sample was also representative of the soil in the area and confirmed its homogeneity. Some small spectral differences were observed with the < 1 mm fines sample from Station C (14141,37) in the vicinity of Cone Crater. The most noticeable feature was the slight strengthening of the band about 450 cm^{-1} which showed that the Fra Mauro soil in this area was contaminated with the more pyroxene-rich Cone Crater soil ejecta. Some rock fragments from the interior of 14321 (associated with piece 14321,263) again showed some spectral enhancement of the features but they were not as pronounced as 14321,98 which was a polished butt end (presumably from the exterior). Again it would appear in fact that the interior pieces of this rock investigated more closely resembled the surrounding soil. In Fig. 6 the Apollo 15 soils show a larger variety of spectral differences that can be associated with their geographic locations and the corresponding mineralogical compositions. The two samples from the Apennine Front at Station 2 (15091,65 and 15221,71) comprise St. George Crater ejecta and other ejecta materials. The two showed some very slight spectral variations but in fact more closely resembled the Apollo 14 soils. Consequently, it would appear that the soils in the vicinity of the St. George Crater rim are substantially intermixed. To some extent this was also true of 15071,60 from Station 1 which might be expected to be more representative of Elbow Crater ejecta. Evidently, this is not the case and again here the soil is well intermixed. In the vicinity of the LM on the mare surface of Palus Putredinis are the faint rays from either of the craters Aristillus and Autolycus. The spectrum of the contingency sample 15021,159 is presumably representative of the fine-grain regolith and although the spectral features were not significantly different there was a lowering of some frequencies which is usually indicative of an increasing iron content and a corresponding slight decrease in the silica and alumina content. The surface layer here may be more representative of Imbrium or other ejecta from deeper crust levels

that would not have traveled as far to the Apollo 15 site compared to the Apollo 14 site. The spectrum of the Dune Crater sample 15471,68 from Station 4 shows further enhancement of the spectral feature at $\sim 450 \text{ cm}^{-1}$ which can be associated with a slight increase in the pyroxene content and this trend is further enhanced, as seen in Fig. 6, in 15531,63 from the rim of Hadley Rille. The spectra are more similar to some of the Apollo 12 rocks such as 12002,186 or 12009,48 which have high pyroxene contents. It may be expected that these young ejecta materials are more representative of bedrock whereas the material from the Apennine Front contains largely gardened mixtures.

The 15426,2 breccia from Station 7 at Spur Crater does not resemble the Fra Mauro breccias as it is obviously richer in pyroxene but it also does not show the same spectroscopic features as the Apollo 12 samples. In Fig. 4 it can be compared with the microbreccia 12073,42 and it also bears some resemblance to the Cone Crater sample 14321,263 in Fig. 5. However, we have not observed any exact spectroscopic analog from any of the soil samples.

SUMMARY

In this paper we have shown that the infrared and Raman spectra provide all the vibrational frequencies of the materials. These comprise the internal stretching and bending modes associated with the $(\text{SiO}_n)^{x-}$ groups, the weak cation-oxygen coupling between adjacent groups and the low-frequency vibrations that can be associated with translational and librational lattice modes. These groups may be isolated, as for example in the olivines, in chains or framework structures which are associated with the pyroxene and plagioclase classes, respectively. Generally as the structure changes from one that is more open to one that has considerable cross coupling, the frequencies of the vibrational modes increase.

The infrared measurements have the advantage that they provide information on the macroscopic properties that may be expected to be more representative of large areas. In consequence, the spectra are more difficult to interpret uniquely but they do indicate spectral differences. The complementary Raman techniques offer a non-destructive tool for determining mineralogical compositions of grains and inclusions and usually provide positively identifiable results.

The Christianson frequencies can be obtained concurrently with the infrared vibrational frequencies from a single infrared reflectivity measurement and a Kramers-Kronig analysis of the spectrum. This frequency, corresponding to the peak in the emittance spectrum, has been shown by others and on a restricted basis by ourselves to be a useful feature for diagnostic applications.

However, we believe that it is necessary to obtain the complete infrared spectrum from $100\text{--}2000 \text{ cm}^{-1}$ in order to correlate the molecular vibrational frequencies and the Christianson frequency with the chemical constituents and to then deduce the major mineral abundances (Perry and Lowndes, 1970). This is of course done by comparison with terrestrial counterparts and an excellent background for this work has been provided by Estep *et al.* (1971 and 1972). Their comprehensive measurements on a large number of terrestrial materials are an important part of the program as the

uniqueness of the vibrational frequencies applied to mineralogical studies has been relatively unexploited by the geologist.

It is apparent from our data that differences in the chemical and mineralogical compositions of the lunar samples can be spectroscopically distinguished. A spectral difference can also be observed between the samples gathered from the different geographical locations, e.g., the spectra of lunar highland material indicate that they are less basic than the Fra Mauro breccias. These materials in turn are less basic than some of the crater ejecta, the Hadley Rille samples, and some of the mare materials from Apollo 11 and 12. These results in consequence have a direct bearing on the origin of the samples and comparison of spectra from various locations (either obtained from the Apollo missions or using remote sensing techniques) will be directly applicable as a petrological tool and will contribute to unraveling the history of the lunar surface.

The reflectance data provide the infrared dielectric dispersion (ϵ' , ϵ'') and the optical constants, n , the refractive index and α , the absorption coefficient, over the range 10–5000 cm^{-1} . These properties do not depend on the type of measurement and they provide fundamental parameters for synthesizing spectra (e.g., emission spectra) for spectral matching models. Another product of the Kramers-Kronig analysis are values of the far-infrared dielectric constant at $\sim 10^{11}$ Hz. These results are less sensitive to sample dryness compared with the low-frequency dielectric constant obtained from capacitance techniques. The temperature dependence and pressure (density) dependence of the dielectric constants extrapolated to lower frequencies have important applications in radar studies of the depth and composition of the various surface layers. Our far-infrared and dielectric studies on lunar basalts are in agreement, but the difference in the results for the samples containing high abundances of iron and titanium indicate other relaxation processes are present at lower frequencies.

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