

The crystal structure of paragonite-2M₁

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Abstract

The structure of a near end member paragonite-2M₁ was refined in C2/c symmetry to a residual of 4.5% with 698 non-zero independent reflections. Attempted refinement in Cc symmetry was not successful. Tetrahedral Si,Al cations are disordered, and the mean T–O distances are in accord with prediction. The small interlayer Na cation facilitates a larger tetrahedral rotation angle (16.2°), smaller *b*-repeat (8.898Å), and smaller interlayer separation (3.053Å) than observed for muscovite-2M₁. It also allows an offset of adjacent 2:1 layers across the interlayer gap that significantly reduces the β angle and is due to mutual repulsion between oxygen anions brought into proximity by tetrahedral tilting.

Introduction

Paragonite, the Na analogue of muscovite, is an abundant rock-forming mineral in low-grade and medium-grade metamorphic rocks. There is a solid solution between paragonite and muscovite and a slightly asymmetric unmixing solvus. Naturally occurring end member paragonite has not been reported.

The most abundant polytype for paragonite is 2M₁, although the 3T form also has been recognized in nature and the 1M form can be made synthetically. The 2M₁ and 3T paragonite polytypes may coexist with very similar compositions (Nakamura and Kihara, 1977). Burnham and Radoslovich (1964) refined the structures of coexisting natural sodian muscovite-2M₁ (K_{0.65}Na_{0.35}) and potassium paragonite-2M₁(Na_{0.85}K_{0.15}) by single crystal X-ray techniques. Soboleva *et al.* (1976) used high voltage texture electron diffraction patterns to refine the structure of a fine grained synthetic paragonite-1M (Na_{0.91}K_{0.09}). Sidorenko *et al.* (1977a,b) applied the same technique for the refinement of natural fine grained 2M₁ (Na_{0.06}K_{0.10}Ca_{0.03}□_{0.27}) and 3T (Na_{0.71}K_{0.16}Ca_{0.03}□_{0.10}) forms of paragonite. This paper reports refinement of a near end member paragonite-2M₁ (Na_{0.92}K_{0.04}Ca_{0.02}□_{0.02}) structure by single crystal X-ray methods. The research was investigated in order to provide more accurate parameters of a paragonite without appreciable interlayer K or vacancies to serve as a reference for comparison with the structures of other Na-micas. Attention will be focused on the interlayer region.

Experimental

Several homogeneous paragonites of near end member composition (~4 mole% muscovite) were kindly furnished by Prof. Dr. Martin Frey of the University of Basel and Professor C. V. Guidotti of the University of Maine. A suitable but thin lath-shaped single crystal 0.32 × 0.25 × 0.018 mm in size was obtained from sample PVB 1705, originally collected by Prof. Dr. Peter Bearth of the University of Basel from a glaucophane-bearing metamorphosed eclogite from the ophiolite zone of Zermatt-Saas Fee in the Swiss Alps. The occurrence is described by Bearth (1973) and Bearth and Stern (1979). An electron microprobe analysis (Table 1) of this sample by C. A. Geiger of this department shows that only one mica is present and gives a resultant formula of (Na_{0.916}K_{0.042}Ca_{0.018}□_{0.024})(Al_{1.990}Fe_{0.028}Mg_{0.013}Ti_{0.003})_{2.034}(Si_{2.939}Al_{1.061})O₁₀(OH)₂.

Cell dimensions of $a = 5.128(2)\text{\AA}$, $b = 8.898(3)\text{\AA}$, $c = 19.287(9)\text{\AA}$, and $\beta = 94.35(3)^\circ$ were obtained by least squares refinement of 15 reflections on a Nicolet P2₁ automated single crystal diffractometer. These also indicate a near end member composition. Data for 5449 reflections were obtained out to $2\theta = 60^\circ$ in all eight octants of the limiting sphere. The $2\theta:\theta$ variable scan technique with monochromatic MoK α radiation was used. Crystal and electronic stability were checked after every 50 reflections by monitoring a standard reflection. Integrated intensities were corrected for Lp and absorption effects and symmetry-averaged to 698 non-zero independent monoclinic reflections. The absorption correction was produced by the empirical ψ scan technique in which the data are compared to complete ψ scans (10° increments of ϕ) for selected reflections spaced at 2θ

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Table 1 Electron microprobe analysis (wt.%).

	1	2	3	4	mean	atoms		
SiO ₂	46.10	45.90	45.36	46.45	45.953	Si	2.939	} 4.00
						Al ^{IV}	1.061	
Al ₂ O ₃	40.51	40.53	40.79	40.06	40.473	Al ^{VI}	1.990	} 2.034
TiO ₂	0.05	0.06	0.06	0.12	0.073	Ti	0.003	
FeO	0.56	0.50	0.49	0.58	0.533	Fe ²⁺	0.028	} 0.976
MgO	0.15	0.12	0.11	0.15	0.133	Mg	0.013	
MnO	0.01	0.00	0.00	0.00	0.003	Mn	—	} 0.976
CaO	0.24	0.29	0.29	0.25	0.267	Ca	0.018	
BaO	0.03	0.00	0.00	0.02	0.013	Ba	—	} 0.976
K ₂ O	0.63	0.45	0.45	0.52	0.513	K	0.042	
Na ₂ O	7.36	7.39	7.34	7.46	7.388	Na	0.916	} 0.976
H ₂ O (diff.)	4.36	4.76	5.11	4.39	4.651			

intervals of 5°. There were no violations of the C-centering or *c*-glide extinctions.

Refinement

Least squares refinement with program ORFLS was initiated in *C2/c*, the ideal space group of the 2M₁ structure. Scattering factors appropriate for 50% ionization were taken from Cromer and Mann (1968). Unit weights were used throughout refinement because experience in this laboratory has shown them to be superior to sigma weights for the intensity distribution found in layer silicates, where there is a great differential in intensity between $k = 3n$ and $k \neq 3n$ indices. Atomic coordinates for the structure of potassian paragonite (Burnham and Radoslovich, 1964), kindly furnished by C. W. Burnham, were used as the initial coordinates for several cycles of refinement. Fixed isotropic B values were used at first and no cation ordering was assumed in the choice of scattering factors. Variable isotropic B values were used, followed by variable anisotropic values. Refinement proceeded smoothly from an initial residual of 0.157 to a final *R* of 0.045. On a final difference map, all atom positions were close to zero density (range 0.02 to 0.06 electrons/Å³) and the position of the proton of the OH group was visible as a peak of 0.24 electrons/Å³. The position of the hydrogen was not refined. Tables 2–6 report the experimental data, final atomic coordinates, thermal ellipsoid orientations, bond lengths, and other important structural features.²

Because refinement in the ideal space group *C2/c* indicated no ordering of the tetrahedral cations, refine-

ment in subgroup *Cc* also was attempted. The procedure used for margarite-2M₁ and muscovite-2M₁ by Guggenheim and Bailey (1975) was followed. Only one ordering pattern of Si and Al in tetrahedral sites is possible in *Cc* symmetry that is consistent with Si,Al disorder in sites T(1) and T(2) of the *C2/c* space group. This is the ordering pattern found in margarite-2M₁. Non-centric atomic coordinates for this pattern were derived by use of distance-least-squares program OPTDIS, and refinement of these coordinates was attempted with program ORFLS. One octahedral cation was held constant in position in order to fix the origin in this noncentrosymmetric *Cc* space group. This refinement was not successful. Although the Al-rich tetrahedra remained larger than the Si-rich tetrahedra, individual T–O bond lengths differed by as much as 0.25 Å within the same tetrahedron and convergence ceased near *R* = 10%.

Discussion of the structure

Many of the structural features of paragonite-2M₁ are very similar to those found in other dioctahedral micas and need not be discussed further. These features include the shapes of the individual tetrahedra and octahedra, the tetrahedral and octahedral sheet thicknesses, and the mean T–O and M–O,OH distances (Tables 5 and 6).

Tetrahedral positions and bond lengths

The angle of tetrahedral rotation of 16.2° (Table 6) is several degrees larger than reported for muscovite (11.3° in the structure by Rothbauer (1971)). This is correlated with the observations that the *b* repeat of paragonite-2M₁ (8.898 Å) is smaller than that of muscovite-2M₁ (9.105 Å) and that the interlayer separation is smaller (3.053 Å vs. 3.480 Å). Radoslovich (1963) pointed out that in muscovite the interlayer K is too large for its 6-coordinated hole after full tetrahedral rotation has taken place to relieve the lateral misfit between the tetrahedral and octahedral sheets. As a result the interlayer K prevents rotation from proceeding to its full extent; the octahedral sheet is stretched laterally to compensate for the incomplete rotation, and adjacent 2:1 layers are propped apart by 0.8 Å by the K relative to pyrophyllite. Verification of this effect was provided by Burns and White (1963) who showed that the *b* repeat distance decreases as K is leached from muscovite. Na is 0.36 Å smaller in ionic radius than K, and allows both greater tetrahedral rotation to relieve the lateral misfit and a closer approach of adjacent layers. The basal oxygen surfaces are still not in contact, however, as judged by the $(c \sin \beta)/2$ value of 9.62 Å in paragonite-2M₁ relative to 9.20 Å in pyrophyllite-I_{Tc} where there is no interlayer cation and the surfaces should be in contact.

The tetrahedral rotation value of 16.2°, the mean Na–O value of 2.62 Å, and the interlayer separation of 3.053 Å are quite close to the values given by Burnham and Radoslovich (1964) for potassian paragonite (15.9°,

² Table 2 of observed and calculated structure amplitudes and Table 4 of thermal ellipsoid orientations may be obtained by ordering document AM-83-236 from the Business Office, Mineralogical Society of America, 2000 Florida Avenue, N.W., Washington, D.C. 20009. Please remit \$5.00 in advance for the microfiche.

Table 6. Other structural features of paragonite-2M₁

Tetrahedral rotation	α		16.2°
Tetrahedral angle	τ	T(1)	110.4°
		T(2)	110.3°
Octahedral rotation	ω		8.3°
Octahedral flattening	ψ	M(1)	62.1°
		M(2)	57.0°
		Na	54.4°
Octahedral distortion*	RMS		($\sigma\theta$) ²
		M(1)	7.24 107.7°
		M(2)	5.43 58.9
Sheet thickness:	tetrahedral		2.243Å
	octahedral		2.078Å
Interlayer separation			3.053Å
Basal oxygen corrugation Δz			0.226Å
Intralayer shift			+ 0.374 _{a2,3}
Layer offset			+ 0.045 _{a1,1}
Resultant shift			- 0.285 _{a1}

*RMS octahedral distortion parameter is defined by Dollase (1969) as the rms-deviation of the 15 octahedral angles from their ideal values. $(\sigma\theta)^2$ is defined by Robinson *et al.* (1971) as $\frac{1}{15} \sum_{i=1}^{15} (\theta_i - 90^\circ)^2 / 11$, where θ_i is the observed O-H-O angle.

from 2.56Å to 2.62Å, and the interlayer separations from 3.00 to 3.06Å. The predicted rotation angle for the paragonite of this chemical composition is 13.7° according to the equation of McCauley and Newnham (1971) and using their suggested bond length values. The difference between the observed and predicted angles suggests that the latter study needs revision to incorporate newly refined mica structure data and new values for bond lengths.

The mean T-O bond lengths (Table 5) for T(1) and T(2) are nearly identical; hence, tetrahedral Si and Al show long-range disorder in paragonite-2M₁, just as in muscovite-2M₁. The observed mean T-O values of 1.652Å and 1.653Å are in excellent agreement with the value of 1.652Å predicted by the regression equation of Hazen and Burnham (1973) for micas. Mean T-O values for the other four paragonite structures reported in the literature are 1.666Å (2M₁ structure by Sidorenko *et al.* (1977a)), 1.652Å (2M₁ by Burnham and Radoslovich (1964)), 1.659Å (1M by Soboleva *et al.* (1976)), and 1.646Å (3T by Sidorenko *et al.* (1977b)). Drits (1975) noted that T-O bond lengths tend to be slightly greater for trioctahedral layer silicates than for dioctahedral forms, but Zvyagin and Soboleva (1979) believed that the paragonites were exceptions to this generalization. They explained that the larger T-O bonds observed in the paragonites studied by electron diffraction were a result of attraction of the basal oxygens by the small Na ions imbedded in the ditrigonal cavities. This was considered to cause both a large rotation angle and a lengthening of the T-O bonds. This effect is not substantiated by the present study.

Beta angle and interlayer region

Each 2:1 layer of the 2M₁ structure ideally has a stagger or intralayer shift of $a/3$ of the upper tetrahedral

sheet relative to the lower tetrahedral sheet. Adjacent layers are positioned so that the directions of stagger alternate regularly by angles of +120°, -120°, +120° If the hexagonal rings of adjacent layers exactly superimpose around the interlayer cation there is a resultant shift of $-a/3$ between layer 1 and layer 3 to give a resultant beta angle of 95°. In most K-micas the observed resultant shift and β angle are larger than ideal as a consequence of octahedral flattening and rotation and of M(1) being larger than M(2). These structural effects cause the stagger within each layer to be appreciably larger than $a/3$, and in addition there is a small offset of adjacent layers around the interlayer K. In the paragonite of this study the intralayer staggers are +0.374 a instead of 0.333 a , quite similar to those in the 2M₁ K-micas, but adjacent 2:1 layers are offset around the interlayer Na by +0.045 a_1 = 0.23Å. This offset is five times the magnitude of the offset reported for muscovite-2M₁ by Rothbauer (1971), and is opposite in direction so that β is reduced in size rather than increased, (β = 94.35° relative to 95.74°).

A similar reduction in β angles has been noted in the other paragonite structures in the literature. Soboleva *et al.* (1977) explain the observed offsets of layers around the interlayer Na as the overall result of mutual repulsion between tetrahedral cations of adjacent layers, interaction of the H proton with the interlayer cation, and mutual repulsion of oxygen anions across the interlayer gap. Although the first factor is strongest it does not provide the required lateral offset, which must be due to the second and third factors. The authors attribute the positive direction of layer offset observed in paragonite-1M to a relatively small corrugation of the basal oxygen surface (Δz = 0.10Å) so that the interaction of the H proton (directed toward the vacant M(1) site) with the interlayer cation becomes dominant in the 1M form. According to this explanation the more corrugated basal surfaces found in most other dioctahedral micas would accentuate repulsion between the closest oxygens of adjacent layers and would tend to displace the upper layer in the opposite direction (along $-a_1$ rather than $+a_1$) relative to the lower layer.

The explanation above requires modification because the H proton has been shown to be quite pliant in position in dioctahedral micas due to its repulsion by the interlayer cation. It is not likely to be able to cause an offset of an entire 2:1 layer by 0.23Å. For paragonite-1M, discussed by Soboleva *et al.* (1977), it is also true that the symmetry precludes any resultant shift of the Na by the H proton, because the 2-fold rotation axis through the Na relates a H⁺ below to an oppositely directed H⁺ at an equal distance above. For the paragonite-2M₁ of this study the H⁺ protons are on the same side of the Na and would tend to repel it along $-b_1$ (Fig. 1) if forceful enough. No such effect is noted. Instead the corrugation of the basal oxygen surfaces (Δz = 0.23Å) due to tetrahedral tilting around the vacant M(1) site places O(3) atoms of adjacent layers in closest contact (3.102Å), and repulsion between

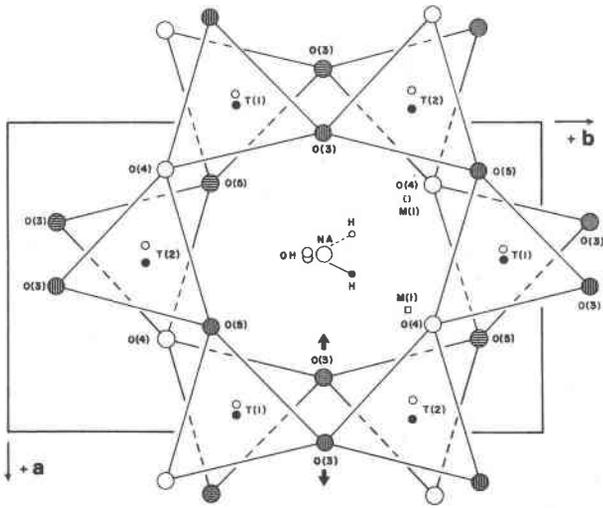


Fig. 1. Projection of interlayer region between layers 1 and 2 onto (001). Tetrahedral tilting around vacant M(1) sites moves O(4) atoms (unshaded) away from median plane at $z = 0.25$. Repulsion between O(3) atoms in closest contact offsets upper layer along $+a$. Cations above median plane shown as solid circles.

them shifts the upper tetrahedral sheet in the $+a_1$ direction of the $2M_1$ cell by an average of $0.045a = 0.23\text{\AA}$ (Fig. 1). For paragonite-1M the resultant offset of layers along $+a_1$ of the 1M cell is smaller both because the different geometry of layer stacking causes repulsion along the $-a_2$ and $-a_3$ pseudohexagonal axes at the same time (with a smaller resultant along $+a_1$) and because adjacent O(3) atoms do not approach as closely (3.196\AA) across the interlayer gap as for paragonite- $2M_1$ (3.102\AA). Although the interlayer geometry is the same for muscovite- $2M_1$ as for paragonite- $2M_1$, the large size of the interlayer K apparently blocks any significant layer offset. The dominance of the interlayer cation size in muscovite- $2M_1$ is illustrated by the variability in direction of the small offsets found in different specimens ($-0.005a_1$ per layer by Rothbauer (1971), $-0.005a_1$ by Güven (1971), but $+0.001a_1$ by Burnham and Radoslovich (1964), $+0.002a_1$ by Gatineau (1963), and $+0.002a_1$ by Birle and Tettenhorst (1968)—all as quoted by Bailey (1980) but taking into account that his values are for total offset per unit rather than per layer). These values are measured on orthogonal projections between the positions of the tetrahedral cations in adjacent layers.

The large decrease in interlayer separation in paragonite- $2M_1$ (3.053\AA) relative to muscovite- $2M_1$ (3.480\AA) reduces the normally thick trigonal antiprism around K to an octahedron of nearly ideal shape around Na ($\psi = 54.4^\circ$). The mean Na-O bond length of 2.624\AA is still considerably larger than predicted (2.38\AA) from the effective ionic radii of Shannon (1976). The same is true of the observed mean Na-O bonds in the other paragonite structures (range 2.56\AA to 2.64\AA). The deviation between

observed and predicted values is due partly to the corrugation of the basal oxygen surface that creates an irregular interlayer cavity. The O(3) atoms that approach each other most closely across the interlayer gap also coordinate most closely with the Na (2.496\AA). The O(4) atoms that have been depressed or elevated away from the mean plane of their neighbors by tetrahedral tilting in turn are farthest from the Na (2.718\AA), whereas the O(5) atoms are intermediate in approach. An irregular interlayer cavity due to the corrugation effect also exists around K in muscovite- $2M_1$, but the differences between individual K-O bond lengths are much smaller ($\Delta = 0.04\text{\AA}$) than in paragonite- $2M_1$ ($\Delta = 0.22\text{\AA}$). The closest observed approach of Na-O(3) = 2.496\AA is in accord with the observation that in micas the bond lengths to the interlayer cation tend to be at least 0.1\AA larger than the sums of the ionic radii for six-fold coordination. This deviation increases with decreasing tetrahedral twist and increasing effective coordination number for the cation.

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