# LETTER

# Rietveld structure refinement of perovskite and post-perovskite phases of NaMgF<sub>3</sub> (Neighborite) at high pressures

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## ABSTRACT

Neighborite (NaMgF<sub>3</sub>) with the perovskite structure, transforms to a post-perovskite (ppv) phase between 27 and 30 GPa. The ppv phase is observed to the highest pressures achieved (56 GPa) at room temperature and transforms to an as yet unknown phase upon heating. Rietveld structure refinement using monochromatic synchrotron X-ray diffraction data provide models for the perovskite and postperovskite structures at high pressure. The refined models at 27(1) GPa indicate some inter-octahedral F-F distances rival the average intra-octahedral distance, which may cause instability in the perovskite structure and drive the transformation to the post-perovskite phase. The ratio of A-site to B-site volume ( $V_A/V_B$ ) in perovskite structured NaMgF<sub>3</sub> (ABX<sub>3</sub>), spans from 5 in the zero-pressure high-temperature cubic perovskite phase to 4 in this high-pressure perovskite phase at 27(1) GPa, matching the  $V_A/V_B$ value in post-perovskite NaMgF<sub>3</sub>. Using Rietveld refinement on post-perovskite structure models, we observe discrepancies in pattern fitting, which may be described in terms of development of sample texture in the diamond-anvil cell, recrystallization, or a change of space group to  $Cmc2_1$ , a non-isomorphic subgroup of Cmcm—the space group describing the structure of CaIrO<sub>3</sub>.

**Keywords:** Perovskite, post-perovskite, pressure, diamond anvil cell, X-ray diffraction, Rietveld modeling, structure

# INTRODUCTION

At the lowermost region of Earth's mantle, the D" layer (Dziewonski and Anderson 1981) is the medium through which heat and chemistry from the core may contact the mantle above. As a result, D" is a complex boundary region (Lay et al. 2004), laterally variable and marked by seismic anomalies. Researchers have long sought to understand the origin and dynamics of this distinct layer (Karato and Karki 2001; Lay et al. 1998). Many high-pressure experiments have investigated MgSiO<sub>3</sub> perovskite for explanations to anomalous seismic phenomena observed in D" and recent experiments have uncovered a post-perovskite structure of MgSiO<sub>3</sub> at pressures and temperatures thought to exist at D" (Murakami et al. 2004; Oganov and Ono 2004).

A perovskite (pv) post-perovskite (ppv) transformation is reported to occur in several ABX<sub>3</sub> materials: MgGeO<sub>3</sub> (Hirose et al. 2005), CaIrO<sub>3</sub> (Hirose and Fujita 2005), and others (Tateno et al. 2006). In addition, the ppv structure is reportedly adopted in A<sub>2</sub>X<sub>3</sub> materials Fe<sub>2</sub>O<sub>3</sub> (Ono and Ohishi 2005), and Mn<sub>2</sub>O<sub>3</sub> (Santillan and Shim 2005). The ppv structure is based on the CaIrO<sub>3</sub> model, which is orthorhombic (*Cmcm*, Z = 4), and rarely adopted among oxides and sulfides at room pressure. The CaIrO<sub>3</sub> structure (Rodi and Babel 1965) may be summarized as alternating layers of BX<sub>6</sub> octahedra and A-site cations normal to the *b* axis within a *C*-centered lattice (Fig. 1a). Octahedra share edges along the *a*  axis and corners along the c axis. The significant anisotropy of a CaIrO<sub>3</sub>-type MgSiO<sub>3</sub> structure may explain the zones of strong seismic anisotropy observed within the D" layer (Garnero 2004; Murakami et al. 2004; Tsuchiya et al. 2004).

The post-perovskite structure is now under intense study (Caracas and Cohen 2005; Mao et al. 2006; Murakami et al. 2005; Shieh et al. 2006) and debate. However, Rietveld structure modeling of ppv-MgSiO<sub>3</sub> is not yet documented. Powder statistics of X-ray diffraction data collected from samples in excess of 120 GPa are often compromised due to small scattering volumes present during experiments. Rietveld refinement of ppv structure will confirm the accuracy of the CaIrO<sub>3</sub> (*Cmcm*) model and the study of analog materials at extreme conditions is an alternative route to this goal. Analogs, such as MgGeO<sub>3</sub> (Duffy et al. 2005; Hirose and Fujita 2005), transform at lower pressures, increasing sample size, peak-to-background discrimination, and overall data quality.

The structure of neighborite (NaMgF<sub>3</sub>) (Chao et al. 1961) was recently studied as an analog material for pv MgSiO<sub>3</sub> (O'Keeffe et al. 1979) using Rietveld structure modeling (Martin et al. 2005; Zhao et al. 1993). Previous work (Liu et al. 2005; Martin et al. 2006) finds a pv-ppv phase transition in NaMgF<sub>3</sub> between 28 and 30 GPa, in agreement with enthalpy calculations (Parise et al. 2004).

Previously we have reported the high-pressure phase transformations of  $NaMgF_3$  (Martin et al. 2006). In the following letter, we report results of Rietveld refinement, detailing the high-pressure perovskite and post-perovskite structures of

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NaMgF<sub>3</sub>. From our analysis we reveal a high-pressure instability in the pv-NaMgF<sub>3</sub> structure and test the CaIrO<sub>3</sub> (*Cmcm*) structure model of ppv-NaMgF<sub>3</sub>.

#### **EXPERIMENTAL METHODS**

Synthetic NaMgF<sub>3</sub> was prepared according to published procedure (Zhao et al. 1993). Fine grains of NaMgF<sub>3</sub> were isolated from an aqueous suspension and dried at 60 °C. The fine powder was loaded into diamond-anvil cells with 350  $\mu$ m anvil culets. Tungsten was used to gasket the sample between the diamond anvils and graphite (1 wt%) was used to absorb the heating laser radiation. The sample was insulated from the anvils during heating using thin layers of NaCl, which also served as an internal pressure standard (Decker 1971; Sata et al. 2002).

Synchrotron X-ray diffraction experiments were conducted at the Advanced Photon Source (APS) at GeoSoilEnviroCARS (GSECARS) 13-ID-D, with  $\lambda = 0.3344(2)$  Å, and at the High Pressure Collaborative Access Team (HPCAT) 16-ID-B, where  $\lambda = 0.4018(2)$  Å. At the European Synchrotron Radiation Facility (ESRF) data were collected at ID-27 with  $\lambda = 0.3738(2)$  Å. Beamline details are described previously for both GSECARS (Shen et al. 2005) and ID-27 (Mezouar et al. 2005). Data at HPCAT and ESRF were collected by a MAR345 imaging plate, while that at GSECARS was collected by a MAR CCD (2048 × 2048). Data were reduced using the program FT2D (Hammersley et al. 1996) and any saturated pixels were excluded from the 360° (chi) intensity integration. The program EXPGUI (Toby 2001) for GSAS (Larson and Von Dreele 2000) was used for Rietveld structure modeling. The background, scale factor, and peak profile function type 3 (GW and LY only) amounted to seven variables.

The data presented in this paper were collected during separate runs and after double-sided infrared laser heating (Schultz et al. 2005). Data from pv-NaMgF<sub>3</sub> were collected after heating to temperatures of 2000  $\pm$  200 K (Benedetti and Loubeyre 2004). To avoid high temperature phase transition in ppv-NaMgF<sub>3</sub> (Martin et al. 2006), this phase was not heated beyond 500  $\pm$  200 °C. Because the Debye-Scherrer rings are round and not elliptical, we believe the variation in intensity around the rings results primarily from preferred orientation of sample powder, rather than from differential stress.

#### RESULTS

Rietveld structure refinement of pv-NaMgF<sub>3</sub> from monochromatic X-ray diffraction data indicate the atomic parameters, MgF<sub>6</sub> tilting angles (Zhao et al. 1993), and relative volumes of the A-site and B-site polyhedra ( $V_A/V_B$  ratio) (Thomas 1998) at 27(1) GPa just before the pv to ppv transition (Fig. 2a, Table 1a). The NaCl pressure media is present as both the B1 and B2 phases (Sata et al. 2002). In addition, we observe weak reflections (downward pointing arrows, Fig. 2a) in diffraction patterns above 14 GPa. The first 3 of these peaks at low angle may be indexed as the 022, 023, and 130 of ppv-NaMgF<sub>3</sub>, which would confirm the phase requires a lower pressure to stabilize under shear stress (Liu et al. 2005). We believe the final 2 reflections at high angle

Post-perovskite structure models of NaMgF<sub>3</sub>



**FIGURE 1.** The orthorhombic post-perovskite structure of  $NaMgF_3$  as (a) *Cmcm* and (b) *Cmc2*<sub>1</sub>.

result from N-phase NaMgF<sub>3</sub> (Martin et al. 2006) rather than from graphite or unknown contamination. Our X-ray diffraction data (Martin et al. 2006) show a complete transformation of pv-NaMgF<sub>3</sub> (*Pbnm*) to ppv-NaMgF<sub>3</sub> above 27(1) GPa.

The X-ray diffraction data from ppv-NaMgF<sub>3</sub> show some evidence for non-uniform texture of the Debye-Scherrer rings and while portions show smooth intensity variation about changing angles of chi, there is evidence of sample recrystallization and/or reorientation close to  $111_{Cmcm}$ . Our Rietveld refinements of ppv-NaMgF<sub>3</sub>, using the CaIrO<sub>3</sub>-type (*Cmcm*) structural model, assign too much intensity to the  $020_{Cmcm}$  and  $200_{Cmcm}$ , while assigning too little to the  $002_{Cmcm}$ , and  $111_{Cmcm}$ . Texture due to preferred orientation of sample powder within the diamond-anvil cell may be modeled using a spherical-harmonic correction implemented in GSAS (Von Dreele 1997). However, the validity of this correction to in-situ high-pressure X-ray diffraction data of orthorhombic materials is debatable and in



NaMgF<sub>3</sub> Post-Perovskite Cmcm & Cmc2<sub>1</sub> trials 55(2) GPa  $\lambda$ =0.3344(1) Å, Obsd. and Diff. Profiles



FIGURE 2. In-situ high pressure Rietveld structure modeling of (a) NaMgF<sub>3</sub> perovskite and (b) post-perovskite NaMgF<sub>3</sub>. Portions of the raw two-dimensional data are inset in each plot to show sample texture. Difference curves of Rietveld structure models are plotted below the calculated (solid) and observed (dotted) background-subtracted X-ray diffraction patterns; (b) in descending order: (1) *Cmcm* with fourth-order spherical harmonic preferred orientation correction (POC), (2) *Cmcm* without POC, (3) *Cmc2*<sub>1</sub>without POC. Weak reflections marked with an arrow are indexed as NaMgF<sub>3</sub> post-perovskite.

practice, when carried out to more than four orders in this case, we find the spherical-harmonic correction will over-fit the data to the detriment of the structural model. While justification of the correction is difficult in this case, the correction may be checked by the texture index, a value summing the magnitude of each texture constant (Von Dreele 1997). Based on previous studies (Sitepu et al. 2005; Von Dreele 1997) and our own trials, we find the *Cmcm* structural model requires a cylindrical fourth-order spherical harmonic preferred orientation correction (cylindrical, 5 parameters) with a texture index of 2.14 (Fig. 2b and Table 1c) to achieve a satisfactory fit to the data.

Because it is not clear whether the intensity integration alone is sufficient to overcome effects of preferred orientation in the

TABLE 1.Structure parameters derived from Rietveld refinement<br/>of data for (a) NaMgF3 perovskite at 27(1) GPa, (b) post-<br/>perovskite structure models of NaMgF3 at 54(2) GPa without<br/>correction for preferred orientation as *Cmcm*, (c) with cor-<br/>rection for preferred orientation as *Cmcm*, and (d) without<br/>correction for preferred orientation as *Cmc2*1

(a) x y z Octahedra 1 Mg-F1 (x2)	a (Å) 4.8904(13) Na (4c) 0.9771(19) 0.0802(13) 0.2500 Filt Mg-F2	b 5.2022(6 Mg (4b) 0.0000 0.5000 0.0000 φ(°) 15.2(8) 2 (x2)	6) <b>Mg</b>	c 7.1403(8) F1 (4c) 0.1019(25) 0.4697(17) 0.2500	) ) 1 <b>V</b>	V (Å <sup>3</sup> ) 181.65(4 F2 (8d) 0.6750( 0.3110(1 0.0619(1 (°) 7.8(7) ⊮ (Å <sup>3</sup> )	4) 17) 10) 12)	Φ(°) 23.2(6) V <sub>A</sub> /V <sub>B</sub> (Å <sup>3</sup> )
1.860(4) Residuals MAR345 (IP	1.886 <b>x2</b> ) 0.1190	(7) )	1.9 <b>wR</b> 0.0	17(8) <b>p</b> 098	8 R 0	8.95(8) <b>R(F²)</b> 9.0588		4.07(5) <b>Dwd</b> 0.5810
(b) x y z	<b>a (Å)</b> 2.7145(14) <b>Na (4c)</b> 0.0000 0.2567(13) 0.2500	<i>b</i> 8.393(5) <b>Mg (4a)</b> 0.0000 0.0000 0.0000	)	c 6.8512(23) F1 (4c) 0.5000 0.4242(15) 0.2500	)	V (Å <sup>3</sup> ) 156.09( F2 (8f) 0.5000 0.1402( 0.0676(	7) 8) 12)	
Mg-F1 x2 1.827(5) Residuals MAR CCD	<b>Mg-F2 x4</b> 1.855(6) <b>χ2</b> 0.0222		V <sub>B</sub> (ų) 8.36(1) wRp 0.0038			V <sub>A</sub> /V <sub>B</sub> 3.66(8) <b>R(F<sup>2</sup>)</b> 0.0325		(Å3) Dwd 0.1120
(c) x y z	(Å) 2.7164(16) <b>Na (4c)</b> 0.0000 0.2422(14) 0.2500	<i>b</i> 8.381(5) <b>Mg (4a)</b> 0.0000 0.0000 0.0000	)	c 6.8487(25) F1 (4c) 0.5000 0.4433(12) 0.2500	)	V (Å <sup>3</sup> ) 155.9(1) <b>F2 (8f)</b> 0.5000 0.1385( 0.0579(9)	) 10) 9)	
Preferred or (2,0,0) -1.1883 Mg-F1 x2 1.777(3) Residuals MAR CCD	ientation cons (2,0,2) 0.3942 Mg-F2 1.830(i x <sup>2</sup> 0.0136	<b>x4</b> 6)	m G (4,0 –0.3 V <sub>B</sub> ( 7.88 wRJ 0.00	SAS. Textur ,0) 3545 Å <sup>3</sup> ) 31(82) o 030	re (4 	Index = 2 4,0,2) 1.6891 ${}_{A}/V_{B}$ (Å <sup>3</sup> ) .94(6) 2(F <sup>2</sup> ) .0196	2.142 )	9. C₁ <sup>mn</sup> as (l,m,n) (4,0,4) -2.1180 <b>Dwd</b> 0.1480
(d) x y z	(Å) 2.7143(9) Na (4a) 0.0000 0.2531(6) 0.0036(7)	<i>b</i> 8.393(3) <b>Mg (4a)</b> 0.0000 0.0129(7 0.2556(7	13) 18)	c 6.846(2) F1 (4a) 0.0000 0.9251(7) 0.0000(7)		V (Å <sup>3</sup> ) 155.99(9 F2 (4a) 0.0000 0.6286( 0.1551(9	9) 10) 9)	<b>F3 (4a)</b> 0.0000 0.6554(8) 0.7902(13)
Mg-F1a 1.900(15) Residuals MAR CCD Note: V <sub>A</sub> /V <sub>B</sub>	Mg-F1 1.752( x2 0.0188 ratio was cal	b M 15) 1.1 w/ 0.1	<b>g-F2</b> 805( <b>Rp</b> 0035	<b>x2 Mg</b> 10) 1.9 <b>R(F</b> 5 0.0 idering, V <sub>A</sub>	-F: 75 <sup>2</sup> ) 32	<b>3</b> (10) 5 (V <sub>UC</sub> /Z)-V	V <sub>B</sub> ( <i>I</i> 8.44 <b>Dwd</b> 0.13	<b>Å</b> 3) V <sub>A</sub> /V <sub>B</sub> (Å3) (1) 3.62(5) Ι 10

data, we have investigated six non-isomorphic subgroups (I) of Cmcm (Hahn 1995) and observe a slightly better fit to diffraction patterns using an alternative structure model of ppv described by lower symmetry space group 36 ( $Cmc2_1$ , Fig. 1b) without a spherical harmonic texture correction. This model was investigated previously to describe a layered sulfide at room pressure (Wada and Onoda 1990), but was later refuted (Kim et al. 1997). This structure model allows extra degrees of freedom, amounting to displacement of B-site cations within further asymmetric  $BX_6$  octahedra. Residuals from Rietveld refinement of  $Cmc2_1$ models are improved (Fig. 2b, Table 1d) and better fits to both  $002_{Cmc21}$  and  $(200)_{Cmc21}$  are achieved, however some intensity of the  $111_{Cmc21}$  as well as  $200_{Cmc21}$  remain partially misfit. Rietveld structure models in the other subgroups,  $C222_1$  for example, improve the fit to 111, but are problematic with data at higher angle. Monoclinic models do not significantly improve the fit to the data nor suggest deviation from orthorhombic symmetry. While the data do not provide irrefutable evidence of a symmetry lower than the Cmcm (CaIrO<sub>3</sub>) structural model for ppv-NaMgF<sub>3</sub>, the distortion and asymmetry found in the  $Cmc2_1$  structure model could perhaps stabilize at extreme conditions. High-pressure studies on other compositions with the ppv structure, may help differentiate effects from preferred orientation.

### DISCUSSION

The MgF<sub>6</sub> tilting angles and  $V_A/V_B$  ratio in the structure of  $pv-NaMgF_3$  (*Pbnm*) at 27(1) GPa reveal an instability, which may drive structure transformation to ppv. The  $V_A/V_B$  ratio, an order parameter for perovskite-perovskite phase transitions, is 5 in cubic  $(Pm\overline{3}m)$  perovskite structures and less than 5, though greater than 4, in perovskite structures where octahedra are tilted (Thomas 1998). At 27(1) GPa, the  $V_A/V_B$  value in pv-NaMgF<sub>3</sub> is 4.07(2), matching the  $V_{\rm A}/V_{\rm B}$  value observed in our Rietveld models of ppv-NaMgF<sub>3</sub>(Cmcm), 3.95(5). This ratio is also equal or close to 4 in the ppv-oxides MgSiO<sub>3</sub>, 4.01, (Murakami et al. 2004), MgGeO<sub>3</sub>, 4.28, (Hirose et al. 2005), CaIrO<sub>3</sub>, 4.19, (Rodi and Babel 1965), and Fe<sub>2</sub>O<sub>3</sub>, 3.89, (Ono and Ohishi 2005). This realization is significant because it suggests minimal relative reconstruction of each site is required to transition from pv to ppv. In addition, it may be possible to predict the pressure of a pv-ppv phase transition in other perovskite (or  $A_2X_3$ ) structured materials by extrapolating the change in  $V_{\rm A}/V_{\rm B}$  at high pressure to a value close to 4. Combining previous data (Liu et al. 2005) with our current results, we find that  $V_A/V_B$  for pv-NaMgF<sub>3</sub> decreases linearly with pressure at a rate of  $1.40(2) \times 10^{-2}$  GPa<sup>-1</sup> between 1 bar and 27(1) GPa.

Examination of the Rietveld structure model of pv-NaMgF<sub>3</sub> at 27(1) GPa reveals that tilting of the perovskite octahedra (Glazer 1972; Liu et al. 2005; Thomas 1998) forces some extraoctahedral F-F distances to rival the average intra-octahedral F-F distance. At which point electrostatic repulsion between select pairs of extra-octahedral fluorine atoms may strain the pv octahedral unit enough to drive the transition to ppv structure. The most notable difference between MgF<sub>6</sub> octahedra in pv and ppv is the amount of strain. Octahedra in ppv-NaMgF<sub>3</sub>, as well as the published CaIrO<sub>3</sub>-type oxides, demonstrate twice the average value of shear strain (via strain matrix, see Zhao et al. 1993) as that observed for pv-NaMgF<sub>3</sub>. This significant strain observed of octahedra in CaIrO<sub>3</sub>-type structures (Rodi and Babel 1965) is likely not favored by the B-site cation and may help explain why so few CaIrO<sub>3</sub>-type structures are stable at room pressure. Similarly, shear strain of octahedra in ppv-NaMgF<sub>3</sub> likely contributes to instability in the structure, which transforms to an unknown orthorhombic phase upon laser heating (Martin et al. 2006).

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